



素粒子現象論研究会2022

2023/03/16

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$S^1/Z_2$  オービフォールド上の  $SU(N)$  ゲージ理論  
における同値類の網羅的分析

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Based on: KT, T. Inagaki, arXiv:2301.12938 [hep-th]





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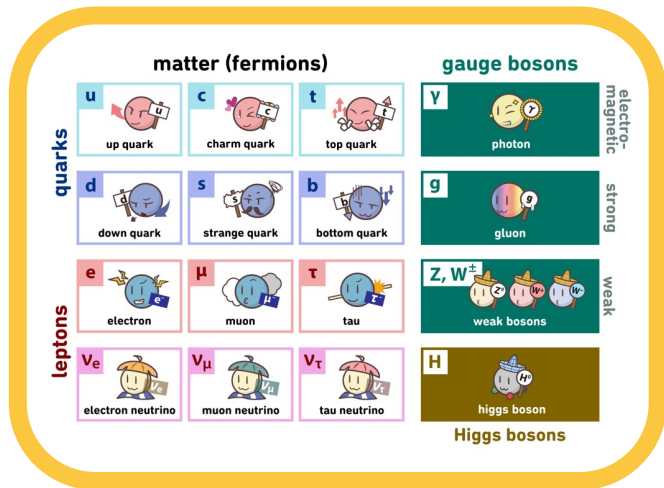
Introduction



# Extra dimensions

[higgstan.com](http://higgstan.com)

Andrew J. Hanson, Indiana University.,  
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SM: 4-dimension

String Theory: 10-dimension





# Extra dimensions

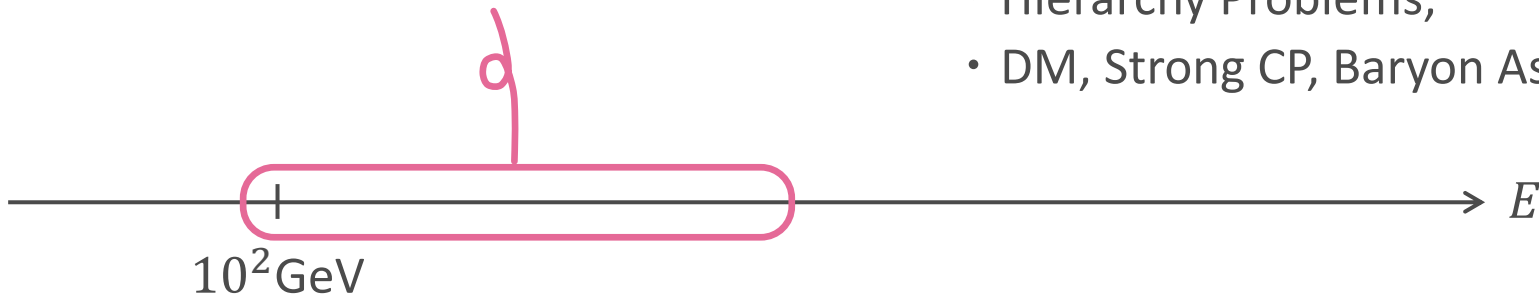
Scenario beyond the SM:  
4-dimension + Extra-Dimension?

## models

- Gauge Higgs Unification (GHU),
- Large Extra Dimension (LED),
- Randall Sundrum (RS), *etc.*

## approach

- Hierarchy Problems,
- DM, Strong CP, Baryon Asym, *etc.*





# Gauge Higgs Unification

Embedding Higgs in extra-dimensional components of gauge fields

$$\underline{x^M = (x^\mu, y) \quad A_M = (A_\mu, A_y) \quad (\mu = 0,1,2,3)}$$

$$\mathcal{L}_{SM}^{4D} = \text{gauge} + \text{fermion} + \text{Higgs} + \text{Yukawa}$$

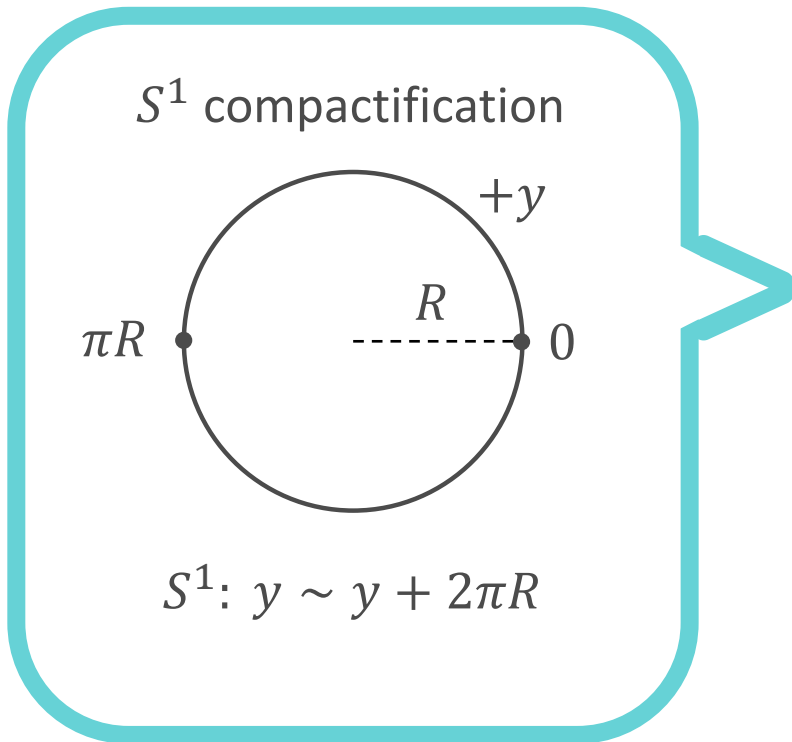


$$\mathcal{L}_{GHU}^{5D} = \text{gauge} + \text{fermion}$$

- Reduced parameters and high predictability
- Naturally solve Hierarchy Problem



# Boundary Conditions (BCs)



infinite interval:  $y \in (-\infty, +\infty)$

$$\phi(y = \pm\infty) = \partial\phi(y = \pm\infty) = 0$$

finite interval:  $y \in [0, 2\pi R)$

$$\phi(y = 0) = +\phi(y = 2\pi R)?$$

$$\phi(y = 0) = -\phi(y = 2\pi R)?$$

$$\phi(y = 0) = e^{i\theta} \phi(y = 2\pi R)?$$

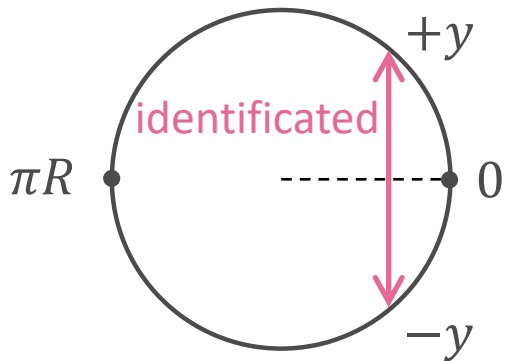
There are many candidates for BCs.



# $S^1/Z_2$ Orbifold

$$\begin{array}{c}
 \xrightarrow{\hat{P}_0} \quad \xrightarrow{\hat{P}_1} \\
 y \quad -y \quad y + 2\pi R \\
 \underbrace{\hspace{10em}}_{\hat{f}}
 \end{array}$$

$S^1/Z_2$  compactification



$$S^1: y \sim y + 2\pi R$$

$$Z_2: y \sim -y$$

$SU(N)$  gauge theory with  $S^1/Z_2$  orbifold

BCs: two parity transformations

$$\begin{cases}
 A_\mu(x, 0 - y) = P_0 A_\mu(x, 0 + y) P_0 \\
 A_\mu(x, \pi R - y) = P_1 A_\mu(x, \pi R + y) P_1
 \end{cases}$$

$P_i$  ( $i = 0, 1$ ) are any  $N \times N$  matrices satisfying,

$$\underline{P_i^2 = 1, \quad P_i^\dagger = P_i}$$

There are many candidates for BCs.





# the Arbitrariness Problem of BCs

Different BCs produce different symmetry.

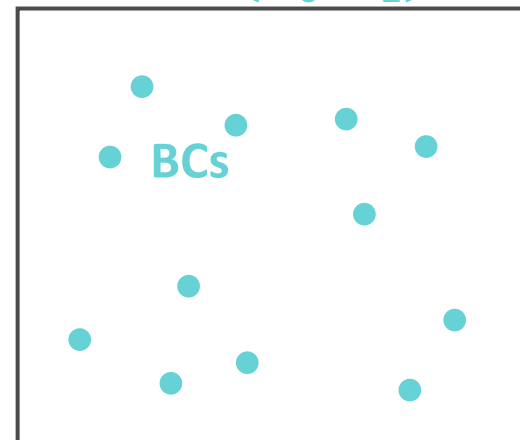
e.g.) SU(3)model

	$P_0$	$P_1$	symmetry
	(+1, +1, +1)	(+1, +1, +1)	$SU(3)$
	(+1, +1, -1)	(+1, +1, -1)	$SU(2) \times U(1)$
diag	(+1, +1, -1)	(+1, -1, +1)	$U(1) \times U(1)$
	(+1, +1, -1)	(+1, $\sigma_2$ )	$U(1)$

Which BCs should be selected? :

the Arbitrariness Problem of BCs

Which  $(P_0, P_1)$ ?



Candidates:



# Equivalence Classes (ECs)

$$P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y),$$

$$\partial_M P'_i = 0, \quad P'_i{}^\dagger = P'_i$$



Hosotani(1989)  
 Haba et al.(2003)  
 Haba et al.(2004)

Some of BCs are connected by a particular gauge transformation:

→ Equivalence Classes (ECs)

$$(P_0, P_1) \sim (P'_0, P'_1)$$

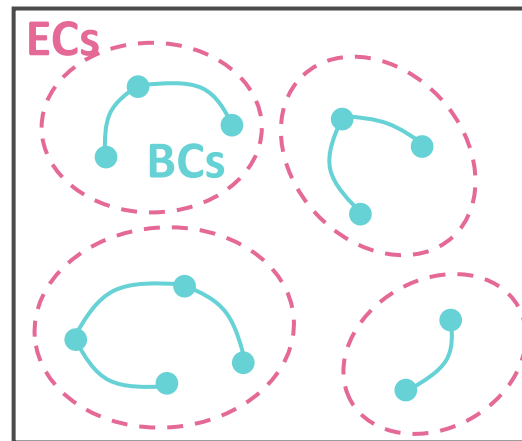
The number of ECs in  $SU(N)$  on  $S^1/Z_2$ :  $(N + 1)^2$

$$N_{+3} C_3 - \sum_{k=0}^{N-2} k_{+1} C_1 N_{-k-1} C_1 = \underline{(N + 1)^2}$$

The remaining arbitrariness:

Which of the  $(N + 1)^2$  ECs should be selected?

Which ECs?



Candidates:



# Our Work

KT, T. Inagaki, arXiv:2301.12938 [hep-th]

$$\partial_M P'_i = 0, \quad P'_i{}^\dagger = P'_i$$

Are there really  $(N + 1)^2$  a lot of ECs?

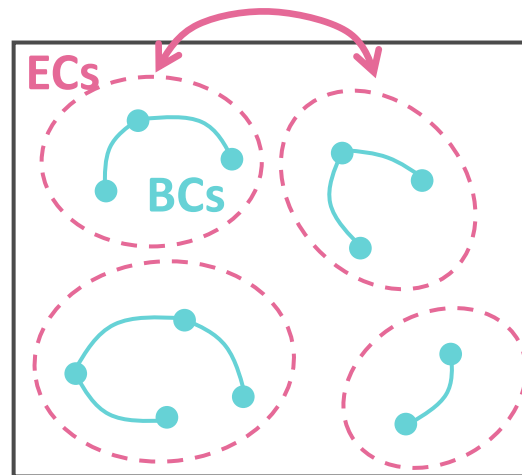
$$(P_0, P_1) \sim (P'_0, P'_1)$$

Previous:

$$P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y), \quad y_0 = 0, \quad y_1 = \pi R$$

$$\text{Gauge transf. } \underline{\Omega(y) = \exp[iy c^a T^a]} \quad \left| \begin{array}{l} T^a: \text{generators} \\ c^a: \text{constants} \end{array} \right.$$

Not connected more?



Candidates:

Analyzing comprehensively by:  $\underline{\Omega(y) = \exp[i f^a(y) T^a]}$

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# Techniques for Analysis

# Simultaneous Diagonalizable

## Tech. 1

Always simultaneous diagonalizable on  $S^1/Z_2$

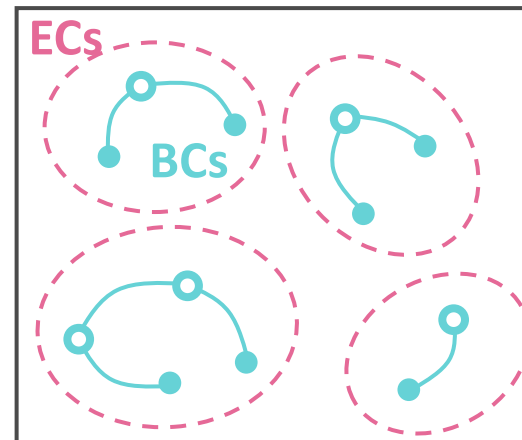
$$(P_0, P_1) \sim (P_0^{\text{diag}}, P_1^{\text{diag}})$$

$$(P_0, P_1)^{\text{diag}} \stackrel{?}{\rightarrow} (P'_0, P'_1)^{\text{diag}}$$

$$P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y),$$

$$\Omega(y) = \exp [i f^a(y) T^a]$$

○: diagonal



Just examine the one  
from diagonal to diagonal



# Classification of Generators

$$y_0 = 0, \quad y_1 = \pi R$$

Kawamura et al.(2020)

Tech. 2

Classify  $U(N)$  generators into 4-types

$$\{T^a\} = \{T^{a++}\} \oplus \{T^{a+-}\} \oplus \{T^{a-+}\} \oplus \{T^{a--}\}$$

Exchangeability with  $P_i$  :

$P_0 \backslash P_1$	commute	anti-commute
commute	$T^{a++}$	$T^{a+-}$
anti-commute	$T^{a-+}$	$T^{a--}$

$$(P_0, P_1)^{\text{diag}} \xrightarrow{?} (P'_0, P'_1)^{\text{diag}}$$

$$P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y),$$

$$\Omega(y) = \exp [i f^a(y) T^a]$$

Requirements for building ECs:

commutative :  $f^a(y_i + y) - f^a(y_i - y) = \text{const.}$

anti-commutative:  $f^a(y_i + y) + f^a(y_i - y) = \text{const.}$

Just examine,  $P'_i = e^{-i(f^a \mp f^a)T^a} P_i$

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Results



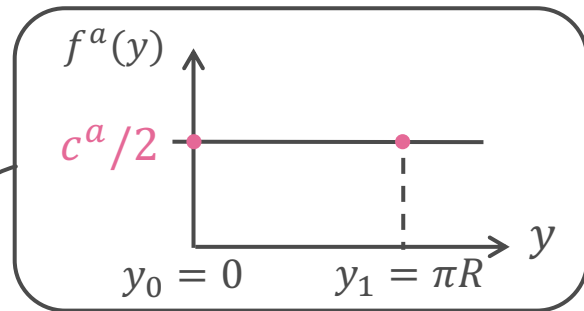
Result:

# for Global parameters

$$\underline{P' = e^{-i(f^a \mp f^a)T^a} P}$$

$$f^a(y) = c^a/2 \quad c^a: \text{constant}$$

$$\left( \begin{array}{l} \text{com. } f^a(y_i + y) - f^a(y_i - y) = 0 \\ \text{anti. } f^a(y_i + y) + f^a(y_i - y) = c^a \end{array} \right) \quad (y_0 = 0, y_1 = \pi R)$$



Generator $\{T^a\}$	$(P'_0, P'_1)$
Commute: $\{T^{a++}\}$	$(P_0, P_1)$
Mixed: $\{T^{a+-}\}$	$(P_0, e^{iT^{+-}} P_1)$
anti-commute: $\{T^{a--}\}$	$(e^{iT^{--}} P_0, e^{iT^{--}} P_1)$

$$T^{+-} = c^{a+-} T^{a+-}$$

$$T^{--} = c^{a--} T^{a--}$$

▷ connect nothing

▷ connect nothing

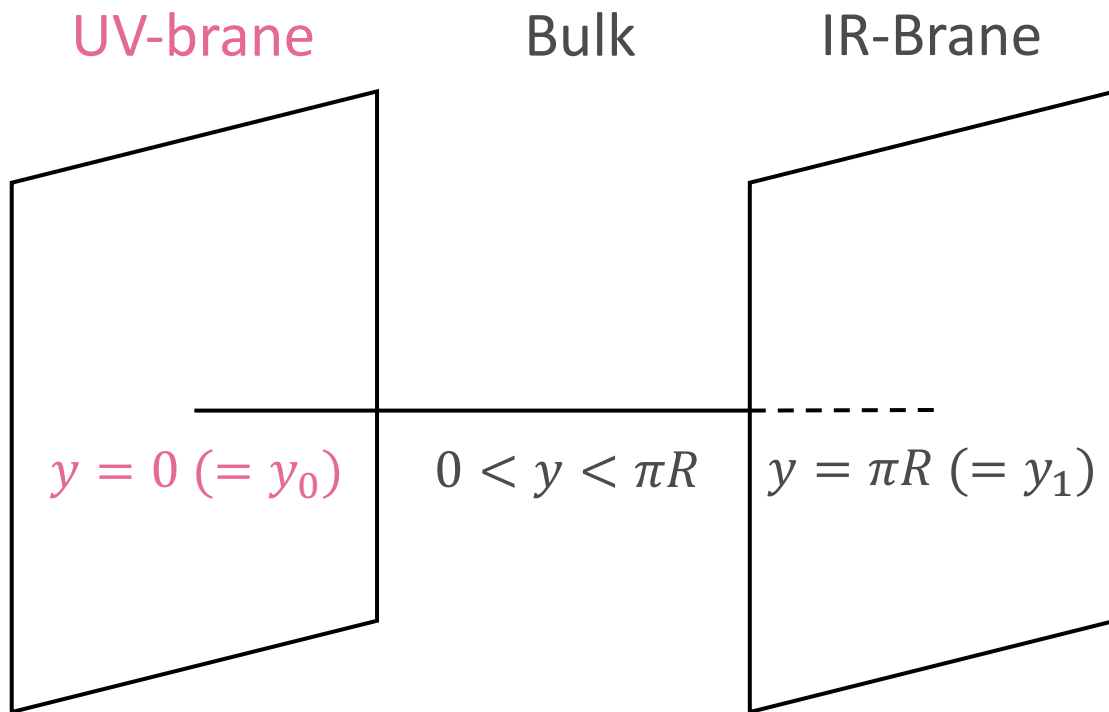
▷ connect nothing





# for Local parameters

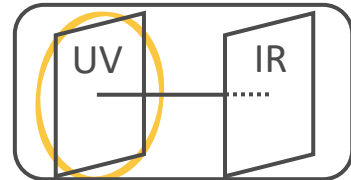
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Result:

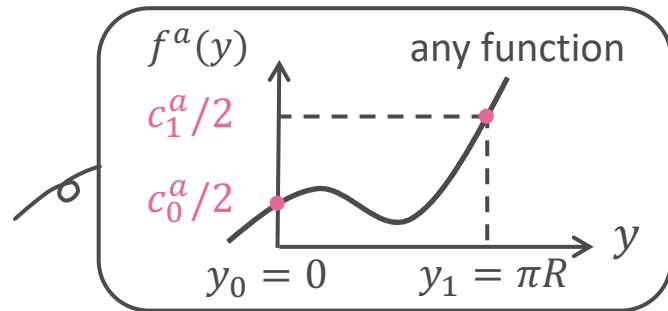
# for Local on UV-brane



$$\underline{P' = e^{-i(f^a \mp f^a)T^a} P}$$

UV limit:  $y \rightarrow 0$

$$\left( \begin{array}{l} \text{com. } f^a(y_i + y) - f^a(y_i - y) \rightarrow 0 \\ \text{anti. } f^a(y_i + y) + f^a(y_i - y) \rightarrow c_i^a \end{array} \right) \quad (y_0 = 0, y_1 = \pi R)$$



Generator $\{T^a\}$	$(P'_0, P'_1)$
Commute: $\{T^{a++}\}$	$(P_0, P_1)$
Mixed: $\{T^{a+-}\}$	$(P_0, e^{iT^{+-}} P_1)$
anti-commute: $\{T^{a--}\}$	$(e^{iT_0^{--}} P_0, e^{iT_1^{--}} P_1)$

$$T_0^{--} = c_0^{a--} T^{a--}$$

$$T_1^{--} = c_1^{a--} T^{a--}$$

▷ connect nothing

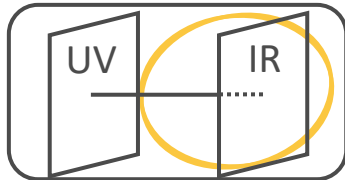
▷ connect nothing

▷ produce  $(N + 1)^2$  18



Result:

# for Local on Bulk & IR-brane

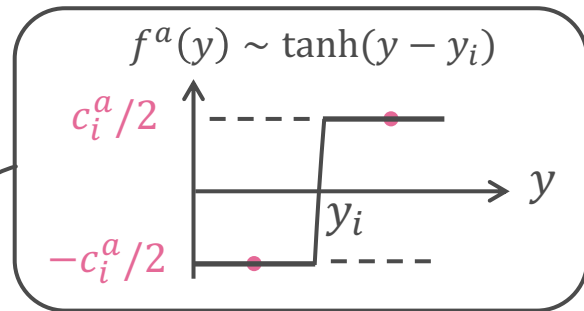


$$P' = e^{-i(f^a \mp f^a)T^a} P$$

Bulk & IR:  $0 < y \leq \pi R$

com.  $f^a(y_i + y) - f^a(y_i - y) = c_i^a (\neq 0)$

achieved by parameters with a kink!



Generator $\{T^a\}$	$(P'_0, P'_1)$
Commute: $\{T^{a++}\}$	$(e^{iT_0^{++}} P_0, P_1), (P_0, e^{iT_1^{++}} P_1)$
Mixed: $\{T^{a+-}\}$	$(P_0, e^{iT^{+-}} P_1)$
anti-commute: $\{T^{a--}\}$	$(e^{iT_0^{--}} P_0, e^{iT_1^{--}} P_1)$

$$T_0^{++} = c_0^{a++} T^{a++}$$

$$T_1^{++} = c_1^{a++} T^{a++}$$



?



connect nothing

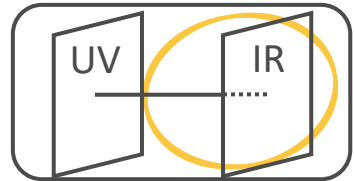


connect  $(N + 1)^2$



Result:

# for Local on Bulk & IR-brane

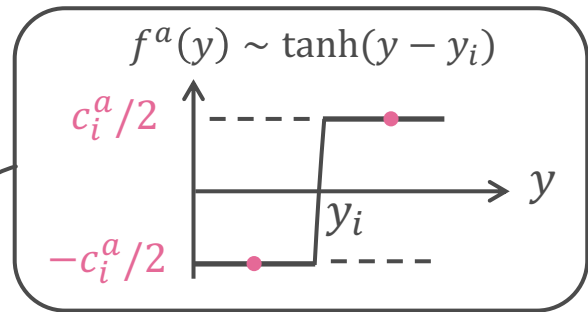


$$P' = e^{-i(f^a \mp f^a)T^a} P$$

Bulk & IR:  $0 < y \leq \pi R$

com.  $f^a(y_i + y) - f^a(y_i - y) = c_i^a (\neq 0)$

achieved by parameters with a kink!



Generator $\{T^a\}$	$(P'_0, P'_1)$
Commute: $\{T^{a++}\}$	$(e^{iT_0^{++}} P_0, P_1), (P_0, e^{iT_1^{++}} P_1)$
Mixed: $\{T^{a+-}\}$	$(P_0, e^{iT^{+-}} P_1)$
anti-commute: $\{T^{a--}\}$	$(e^{iT_0^{--}} P_0, e^{iT_1^{--}} P_1)$

$$T_0^{++} = c_0^{a++} T^{a++}$$

$$T_1^{++} = c_1^{a++} T^{a++}$$

- ▷ connect all set of BCs!
- ▷ connect nothing
- ▷ connect  $(N + 1)^2$



# On the Bulk and IR-brane

$$(P'_0, P'_1) = (P_0, e^{iT_1^{++}} P_1)$$

$$T_1^{++} = \begin{pmatrix} a_1 & & & 0 \\ & a_2 & & \\ & & \ddots & \\ 0 & & & a_N \end{pmatrix}$$

$$\left| \begin{aligned} e^{iT_1^{++}} &= \text{diag}(e^{ia_1}, e^{ia_2}, \dots, e^{ia_N}) \\ &= \tilde{I} \quad \text{: diagonal matrix with } \pm 1 \quad (P'_1{}^\dagger = P'_1) \end{aligned} \right.$$

$$P'_1{}^{\text{diag}} = \tilde{I} P_1{}^{\text{diag}} = \{ \text{All patterns of } P_1 \}$$

 All set of the BCs are connected!

$$P_1 = \begin{pmatrix} + & & & \\ & + & & \\ & & + & \\ & & & + \end{pmatrix}$$



$$P_1 = \begin{pmatrix} + & & & \\ & + & & \\ & & + & \\ & & & - \end{pmatrix}$$

freely flip!

0

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Summary



# Summary

We have comprehensively investigated gauge connected BCs in  $SU(N)$  on  $S^1/Z_2$ .

- There are only  $(N + 1)^2$  Equivalence Classes (ECs) on UV-brane: ( $y = 0$ ).
- The arbitrariness of boundary conditions is completely resolved on Bulk and IR-brane: ( $0 < y \leq \pi R$ ).

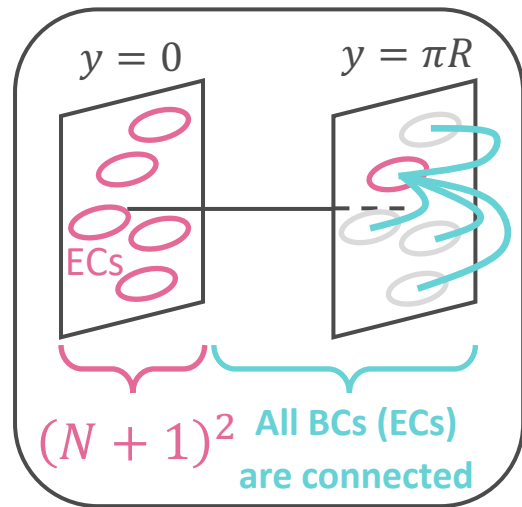
## Future Work

We can freely choose BCs on Bulk & IR

→ Apply to existing  $S^1/Z_2$  models

ECs are unrelated on UV, but related on Bulk & IR

→ Construction of a mechanism to transition between ECs on UV





# Summary

Thank you for listening 

We have comprehensively investigated gauge connected BCs in  $SU(N)$  on  $S^1/Z_2$ .

- There are only  $(N + 1)^2$  Equivalence Classes (ECs) on UV-brane: ( $y = 0$ ).
- The arbitrariness of boundary conditions is completely resolved on Bulk and IR-brane: ( $0 < y \leq \pi R$ ).

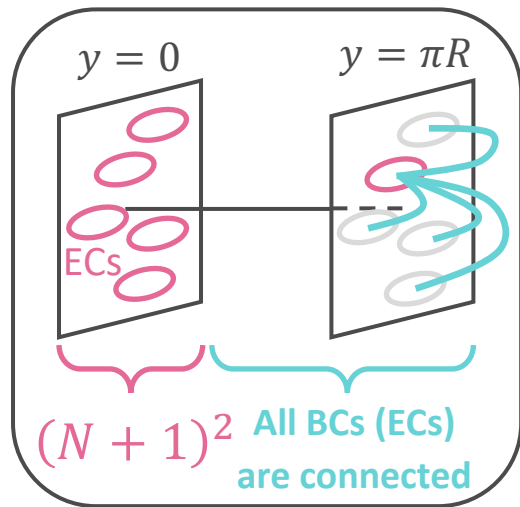
## Future Work

We can freely choose BCs on Bulk & IR

→ Apply to existing  $S^1/Z_2$  models

ECs are unrelated on UV, but related on Bulk & IR

→ Construction of a mechanism to transition between ECs on UV





05

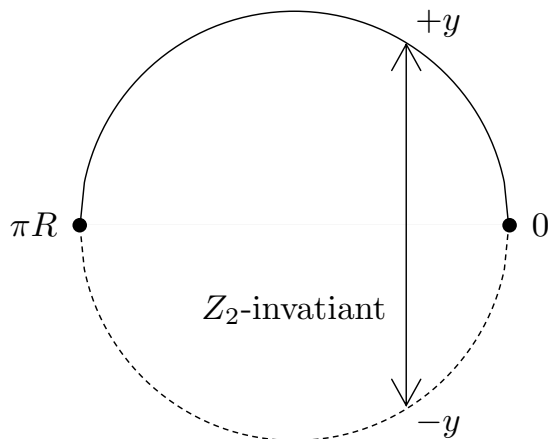
Follow-up



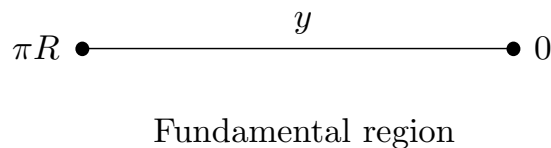
# Two pictures for $S^1/Z_2$ Orbifold

The field are treated in two equivalent pictures on  $S^1/Z_2$ .

(a)  $S^1$  compact space



(b)  $S^1/Z_2$  compact space



Physical fields are defined only on the fundamental region:  $0 \leq y \leq \pi R$



# Step parameters

Parameters with a kink:

$$f_0(y) = \frac{1}{2} \tanh [\lambda (y - (0 - \epsilon))]$$

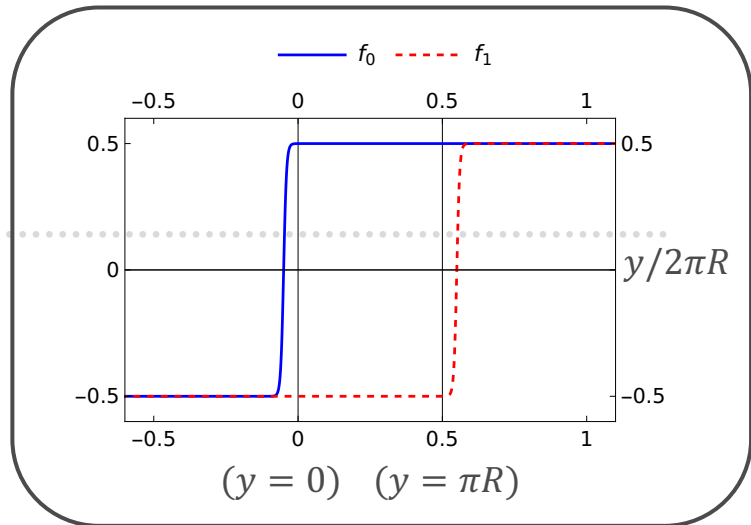
The transformed gauge field and the BCs:

$$A'_y(x, y) = \Omega(x, y) A_y(x, y) \Omega^{-1}(x, y) - \frac{i}{g} \Omega(x, y) \partial_y \Omega^{-1}(x, y)$$

$$A'_y(x, y_i - y) = -P'_i A'_y(x, y_i + y) P_i'^{\dagger} - \frac{i}{g} P'_i (-\partial_y) P_i'^{\dagger}$$

Redefine Bulk & IR-brane area:  $1/\lambda \ll \epsilon < y \leq \pi R$

$\Rightarrow$   $\left| \begin{array}{l} f_0(0 + y) - f_0(0 - y) = 1 \\ f_0(\pi R + y) - f_0(\pi R - y) = 0 \end{array} \right. \rightarrow$  Commutative type produces the phase factor  $e^{iT_{++}}$



Well-defined in the fundamental region



# by Multiple Generators

It is expected that they cannot produce non-trivial diagonal matrix.

generators	c-number components	sub-matrix components
two (ex. $e^{T_{+-} + T_{--}}$ )	Done!	Done!
three (ex. $e^{T_{+-} + T_{-+} + T_{--}}$ )	Done!	Not yet.
full ( $e^{T_{+++} + T_{+-} + T_{-+} + T_{--}}$ )	Not yet.	Not yet.

$$T^{a++} = \begin{pmatrix} \star & 0 & 0 & 0 \\ 0 & \star & 0 & 0 \\ 0 & 0 & \star & 0 \\ 0 & 0 & 0 & \star \end{pmatrix},$$

$$T^{a+-} = \begin{pmatrix} 0 & \star & 0 & 0 \\ \star & 0 & 0 & 0 \\ 0 & 0 & 0 & \star \\ 0 & 0 & \star & 0 \end{pmatrix},$$

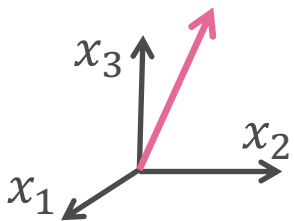
$$T^{a-+} = \begin{pmatrix} 0 & 0 & \star & 0 \\ 0 & 0 & 0 & \star \\ \star & 0 & 0 & 0 \\ 0 & \star & 0 & 0 \end{pmatrix},$$

$$T^{a--} = \begin{pmatrix} 0 & 0 & 0 & \star \\ 0 & 0 & \star & 0 \\ 0 & \star & 0 & 0 \\ \star & 0 & 0 & 0 \end{pmatrix},$$

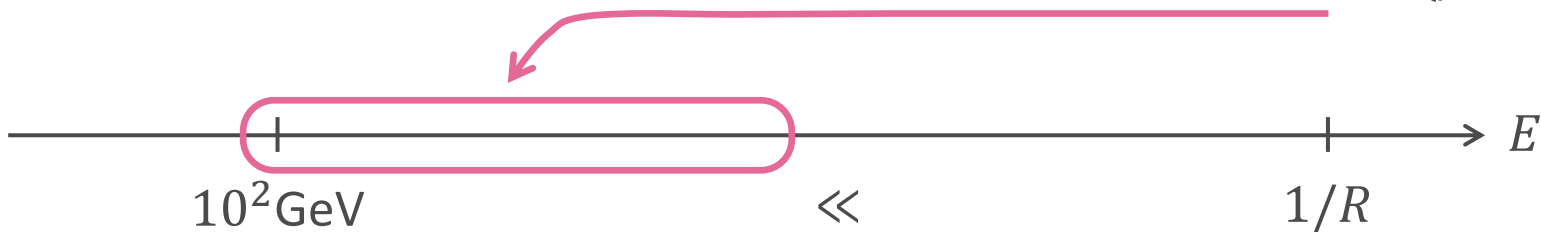
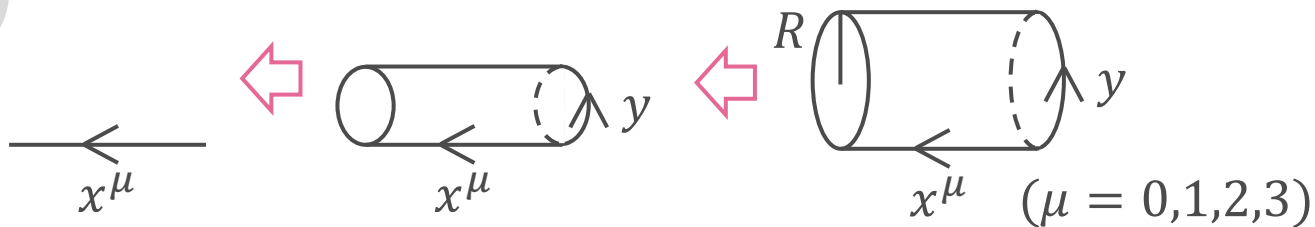


# Compact space

Where??



Extra space is curled up with small radius!





# $SU(N)$ model on $S^1/Z_2$ orbifold

$$\left\{ \begin{array}{l} \mathcal{L} = -\frac{1}{4} F_{MN} F^{MN}(x, y) + \bar{\Psi} i \Gamma^M D_M \Psi(x, y) \\ \Psi(x, y_i - y) = P_i \gamma^5 \Psi(x, y_i + y) \\ A_\mu(x, y_i - y) = P_i A_\mu(x, y_i + y) P_i \quad (i = 0, 1) \\ A_y(x, y_i - y) = -P_i A_y(x, y_i + y) P_i \end{array} \right.$$

$$\left\{ \begin{array}{l} \Gamma^M = (\gamma^\mu, i\gamma^5) \\ \{\Gamma^M, \Gamma^N\} = 2\eta^{MN} I_{4 \times 4} \\ D_M = \partial_M - ig A_M \\ F_{MN} = \frac{i}{g} [D_M, D_N] \\ y_0 = 0, y_1 = \pi R \end{array} \right.$$

Unitary + Parity  $\rightarrow$  Hermite

$$P_i^{-1} = P_i^\dagger \quad P_i^2 = 1 \quad P_i^\dagger = P_i = P_i^{-1}$$

$P_i$  : Hermitian  $N \times N$  matrices  
with  $\pm 1$  eigenvalues



# Symmetry Breaking by BCs

Different BCs generally produce different symmetry.

e.g.) SU(3)model

$P_0, P_1$	$\begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$	$\begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$	$\begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix}$
$A_\mu^0$	$\begin{pmatrix} * & * & 0 \\ * & * & 0 \\ 0 & 0 & * \end{pmatrix}$	$\begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & * \\ 0 & * & 0 \end{pmatrix}$
generators	$\{T^1, T^2, T^3, T^8\}$	$\{T^3, T^8\}$	$\{\frac{\sqrt{3}}{2}T^3 + \frac{1}{2}T^8\}$
symmetry	SU(2) × U(1)	U(1) × U(1)	U(1)



# Gauge transformation for BCs

$$\left\{ \begin{array}{l} \Psi(x, y_i - y) = P_i \gamma^5 \Psi(x, y_i + y) \quad (P_i^\dagger = P_i = P_i^{-1}) \\ A_\mu(x, y_i - y) = P_i A_\mu(x, y_i + y) P_i \\ A_y(x, y_i - y) = -P_i A_y(x, y_i + y) P_i \end{array} \right.$$

↓

$$\left\{ \begin{array}{l} \Psi'(x, y_i - y) = P'_i \gamma^5 \Psi'(x, y_i + y) \\ A'_\mu(x, y_i - y) = P'_i A'_\mu(x, y_i + y) P_i'^{\dagger} - \underline{P'_i \partial_\mu P_i'^{\dagger}} \\ A'_y(x, y_i - y) = -P'_i A'_y(x, y_i + y) P_i'^{\dagger} - \underline{P'_i (-\partial_y) P_i'^{\dagger}} \\ \underline{P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y)} \end{array} \right.$$

Not invariant!





# Gauge transformation for BCs

$$\left\{ \begin{array}{l} \Psi(x, y_i - y) = P_i \gamma^5 \Psi(x, y_i + y) \quad (P_i^\dagger = P_i = P_i^{-1}) \\ A_\mu(x, y_i - y) = P_i A_\mu(x, y_i + y) P_i \\ A_y(x, y_i - y) = -P_i A_y(x, y_i + y) P_i \end{array} \right.$$

↓

$$\left\{ \begin{array}{l} \Psi'(x, y_i - y) = P_i' \gamma^5 \Psi'(x, y_i + y) \\ A'_\mu(x, y_i - y) = P_i' A'_\mu(x, y_i + y) P_i'^\dagger - P_i' \partial_\mu P_i'^\dagger \\ A'_y(x, y_i - y) = -P_i' A'_y(x, y_i + y) P_i'^\dagger - P_i' (-\partial_y) P_i'^\dagger \\ P_i' = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y) = P_i \end{array} \right.$$

$P_i' = P_i \quad (i=0,1)$   
 $\Rightarrow$  gauge invariant



# Gauge transformation for BCs

$$\left\{ \begin{array}{l} \Psi(x, y_i - y) = P_i \gamma^5 \Psi(x, y_i + y) \quad (P_i^\dagger = P_i = P_i^{-1}) \\ A_\mu(x, y_i - y) = P_i A_\mu(x, y_i + y) P_i \\ A_y(x, y_i - y) = -P_i A_y(x, y_i + y) P_i \end{array} \right.$$

↓

$$\left\{ \begin{array}{l} \Psi'(x, y_i - y) = P'_i \gamma^5 \Psi'(x, y_i + y) \\ A'_\mu(x, y_i - y) = P'_i A'_\mu(x, y_i + y) P_i'^{\dagger} - \cancel{\frac{i}{g} P'_i \partial_\mu P_i'^{\dagger}} \\ A'_y(x, y_i - y) = -P'_i A'_y(x, y_i + y) P_i'^{\dagger} - \cancel{\frac{i}{g} P'_i (-\partial_y) P_i'^{\dagger}} \\ P'_i = \Omega(x, y_i - y) P_i \Omega^\dagger(x, y_i + y) \end{array} \right.$$

$$\left\{ \begin{array}{l} \partial_M P'_i = 0 \\ P_i'^{\dagger} = P'_i \end{array} \quad (i=0,1) \right.$$

⇒ not invariant,  
but  $P'_i$  become  
other BCs !



# Physical Symmetry

- Wilson line phase (AB phase)

$$WT \equiv \mathcal{P} \exp \left[ ig \int_0^{2\pi R} dy A_y(x, y) \right] T \quad (T = P_1 P_0)$$

- There is one physics in one EC. Hosotani(1989)

$$V_{eff}(A_M^c; P_0, P_1) = V_{eff}(A_M'^c; P'_0, P'_1)$$

- Physical symmetry

$$(A_y^c; P_0, P_1) \sim (A_y'^c = 0; P_0^{phys}, P_1^{phys})$$

$$\mathcal{H}^{phys} = \left\{ T^a \in \mathcal{G} \mid [T^a, P_i^{phys}] = 0 \ (i = 0, 1) \right\}$$



# ECs in $SU(N)$ on $S^1/Z_2$

- Rearrangement

$$P_0 = \text{diag} \left( \overbrace{+1, \dots, +1, +1, \dots, +1}^N, -1, \dots, -1, -1, \dots, -1 \right),$$
$$P_1 = \text{diag} \left( \underbrace{+1, \dots, +1}_p, \underbrace{-1, \dots, -1}_q, \underbrace{+1, \dots, +1}_r, \underbrace{-1, \dots, -1}_{s=N-p-q-r} \right),$$
$$(P_0, P_1) = [p, q, r, s]$$

- Well-known ECs gauge transformations on  $S^1/Z_2$

$${}_{N+3}C_3 - \sum_{k=0}^{N-2} {}_{k+1}C_1 {}_{N-k-1}C_1 = (N+1)^2$$

$$\begin{aligned} [p, q, r, s] &\sim [p-1, q+1, r+1, s-1] \\ &\sim [p+1, q-1, r-1, s+1] \end{aligned}$$



# Simple Example

e.g.) mixed type for  $U(2)$  matrix

$$\underline{T^{+-} = \begin{pmatrix} 0 & c \\ c^* & 0 \end{pmatrix}}$$

$$e^{iT^{+-}} = \begin{pmatrix} \cos |c| & \frac{ic}{|c|} \sin |c| \\ \frac{ic^*}{|c|} \sin |c| & \cos |c| \end{pmatrix}$$

$$\text{diag}(e^{iT^{+-}}) = I_2, -I_2$$

$$P_0 = \begin{pmatrix} + & \\ & + \end{pmatrix}, \quad P_1 = \begin{pmatrix} + & \\ & - \end{pmatrix}$$



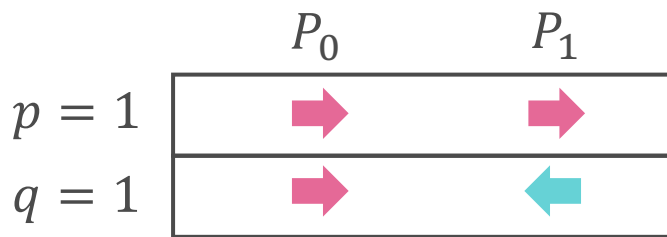
$$P'_0 = P_0, \quad P'_1 = e^{iT^{+-}} P_1$$
$$= \begin{pmatrix} + & \\ & + \end{pmatrix}, \quad = \begin{pmatrix} - & \\ & + \end{pmatrix}$$



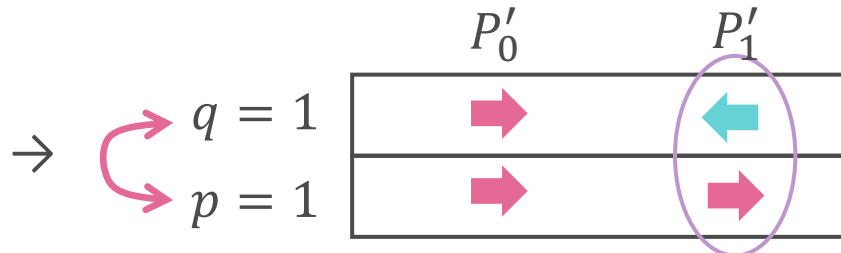
# Simple Example

e.g.) mixed type for  $U(2)$  matrix

$$P_0 = \begin{pmatrix} + & \\ & + \end{pmatrix}, \quad P_1 = \begin{pmatrix} + & \\ & - \end{pmatrix} \quad \Rightarrow \quad P'_0 = \begin{pmatrix} + & \\ & + \end{pmatrix}, \quad P'_1 = \begin{pmatrix} - & \\ & + \end{pmatrix}$$



$[1,1,0,0]$



$[1,1,0,0]$

A set of BCs  $[p, q, r, s]$  are invariant



# General Example

e.g.) anti-commutative type for  $U(N)$  matrix

$$T^- = \begin{pmatrix} 0 & 0 & 0 & A \\ 0 & 0 & B & 0 \\ 0 & B^\dagger & 0 & 0 \\ A^\dagger & 0 & 0 & 0 \end{pmatrix}$$

$A, B$ : sub-matrices

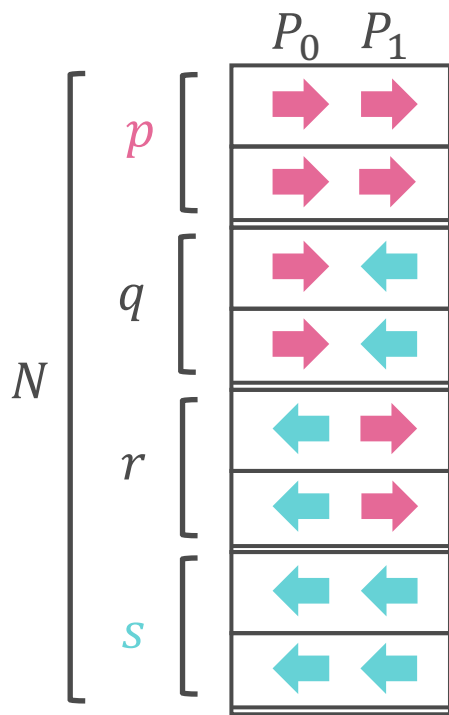
$$\text{diag}(e^{iT^-}) = \begin{pmatrix} \tilde{I}_{p,X} & 0 & 0 & 0 \\ 0 & \tilde{I}_{q,Y} & 0 & 0 \\ 0 & 0 & \tilde{I}_{r,Y} & 0 \\ 0 & 0 & 0 & \tilde{I}_{s,X} \end{pmatrix}$$

$\tilde{I}_{p,X}$ : diagonal ( $p \times p$ ) sub-matrix with  $\pm 1$   
 $X$ : the number of  $-1$  components



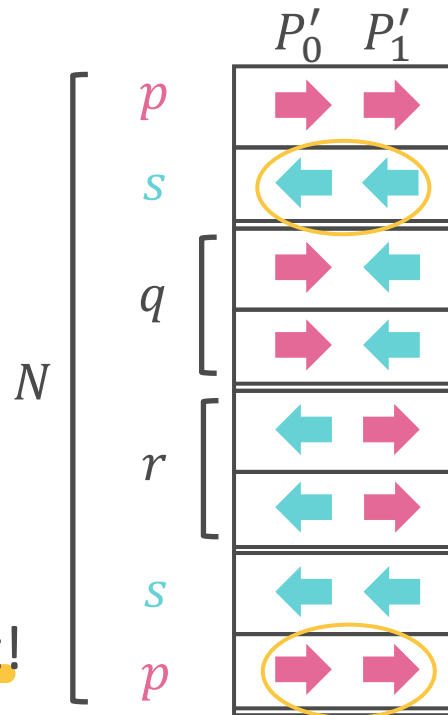
# General Example

e.g.) anti-commutative type for  $U(N)$  matrix



e.g.)  
 $X = 1$   
 $Y = 0$   
 $\rightarrow$

$[p, q, r, s]$  is invariant!





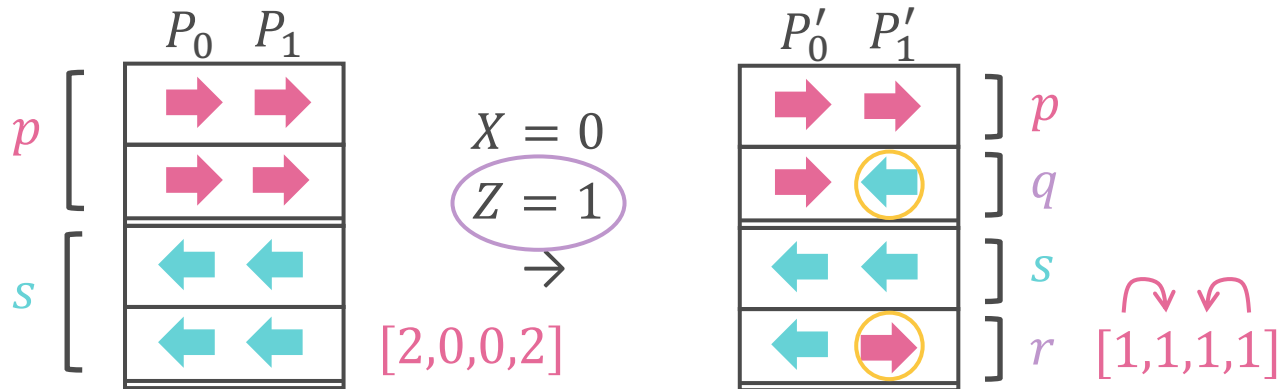
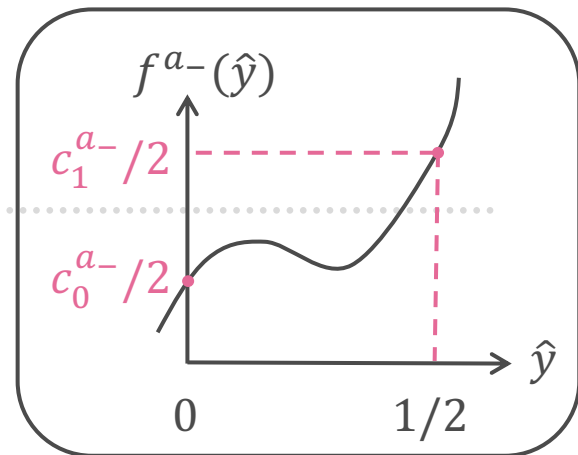


# On the UV-brane

e.g.) anti-commutative type:  $[p, q, r, s] = [2, 0, 0, 2]$

$$\text{diag}(e^{iT_0^-}) = \text{diag}(\tilde{I}_p, \mathbf{X}, \tilde{I}_s, \mathbf{X})$$

$$\text{diag}(e^{iT_1^-}) = \text{diag}(\tilde{I}_p, \mathbf{Z}, \tilde{I}_s, \mathbf{Z})$$

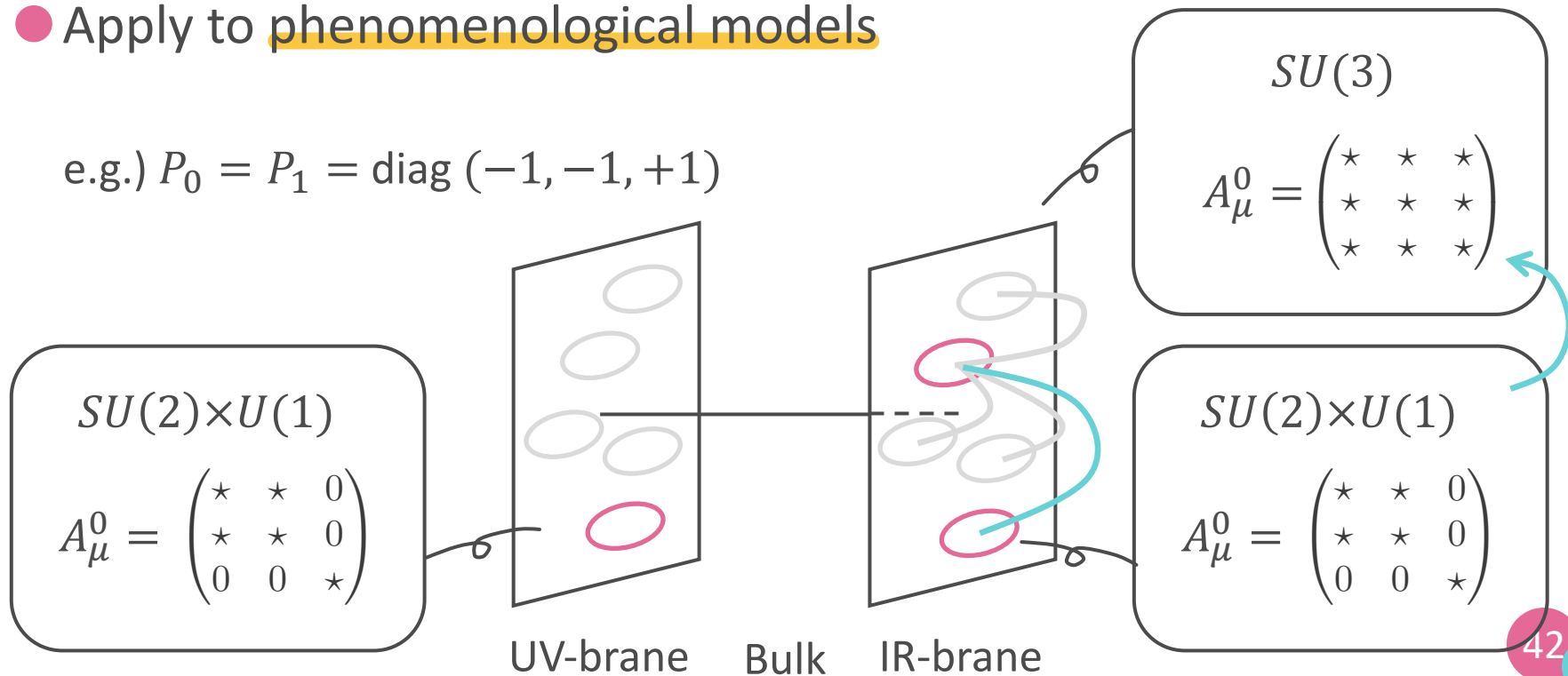




# Applications (1/2)

- Apply to phenomenological models

e.g.)  $P_0 = P_1 = \text{diag}(-1, -1, +1)$



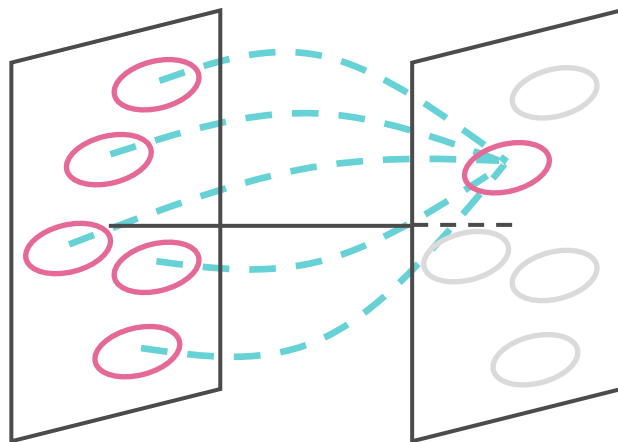


# Applications (2/2)

- Approach to the arbitrariness problem of BCs:

Which type of BCs should be selected without relying on phenomenological information?

Each EC is unrelated



UV-brane

Bulk

IR-brane

All ECs are related