

# Reflection symmetries and universal four-zero texture

素粒子現象論研究会@大阪市立大, Nov. 26, 2020

梁 正樹 (PD) @ 埼玉大学

Arxiv:2003.11701v3 Chin.Phys.Cに投稿中

概要：

Reflection symmetriesというSMの一般化されたCP 対称性を提案

⇒ Majorana位相を  $\alpha_{2,3}/2 = 0$  or  $\pi/2$ に固定 (+  $\theta_{\text{QFD}}^{\text{tree}} = 0$ )

+ universal four-zero texture

⇒ すべてのfermion 質量と  $V_{\text{CKM}}$  &  $U_{\text{MNS}}$  を再現し、

$\delta_{\text{CP}} \doteq 203^\circ$ , 通常質量階層、  $m_1 \doteq 2.5$  or  $6.2$  meVを予言

秋の学会発表

ここまで

+ 対称性の実現のために2HDM+U(1)<sub>PQ</sub> × CPでSSBを考えた

⇒  $f_a \sim M_{\text{GUT}} \sqrt{m_{u,d} m_{c,s}} / v \sim 10^{12} \text{ GeV} \Rightarrow \Omega_a h^2 \sim 0.2$ . axion DM??

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and flavored axion

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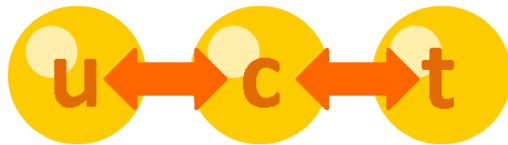
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# Great success of the SM

## 物質粒子

物質を形成する素粒子

クォーク



レプトン



## ゲージ粒子

力を伝える素粒子



## ヒッグス粒子

質量の起源となる素粒子



## Unsolved problems

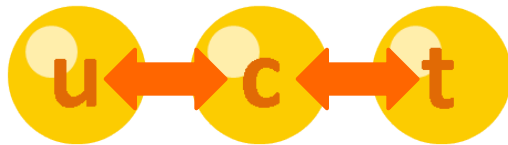
- Neutrino mass
- Flavor puzzle
- Dark matter, etc...

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# 世代の謎 (Flavor puzzle) : 概観

- **Flavor symmetry (top-down)**
  - $U(1)_{FN}$  Froggatt & Nielsen, 1979, **1600**cited
  - $SU(2), SU(3), \dots$
  - $S_3, A_4, S_4, \dots$
  - Family unification Ramond 1979, Wilczek & Zee, 1982  
 $SO(16), SO(18), E_7, E_8, \text{etc}, \supset SO(10) \times SU(3)$
- **Flavor texture (bottom-up)**
  - 6-zero texture Weinberg 1977, Fritzsche 1977, **800**cited
  - $n(=2,3,4..)$ -zero texture
  - Democratic texture Harari, Haut, Weyers, 1978
  - $\mu$ - $\tau$  symmetry Fukuyama and Nishiura 1997
  - Generalized CP sym. Ecker, Grimus, Konetschny '80,  
Holthausen et al, '12, Feruglio, et al, '12
  - Etc..

**世代構造  $\Rightarrow$  Higgsが何かのヒント?**

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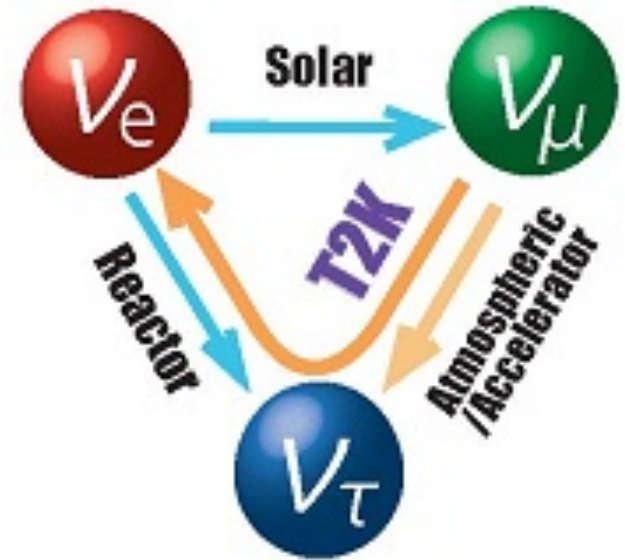
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# Neutrino oscillation

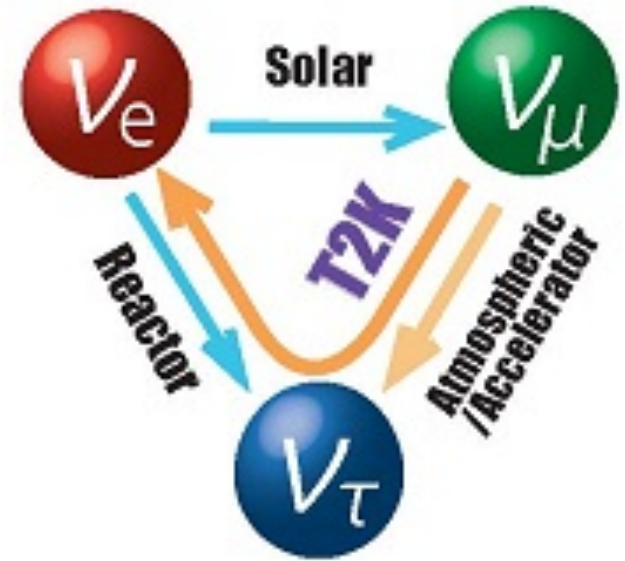
MNS matrix  $U_{\text{MNS}} = U_e U_\nu^\dagger$ .

$$\begin{aligned}
 U &= \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \\
 &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \\
 &\quad \text{atmospheric} \qquad \qquad \text{reactor} \qquad \qquad \text{solar} \qquad \qquad \text{Majorana phase}
 \end{aligned}$$



# Neutrino oscillation

MNS matrix  $U_{\text{MNS}} = U_e U_\nu^\dagger.$



$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \\
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From atmospheric

reactor

solar

Majorana phase

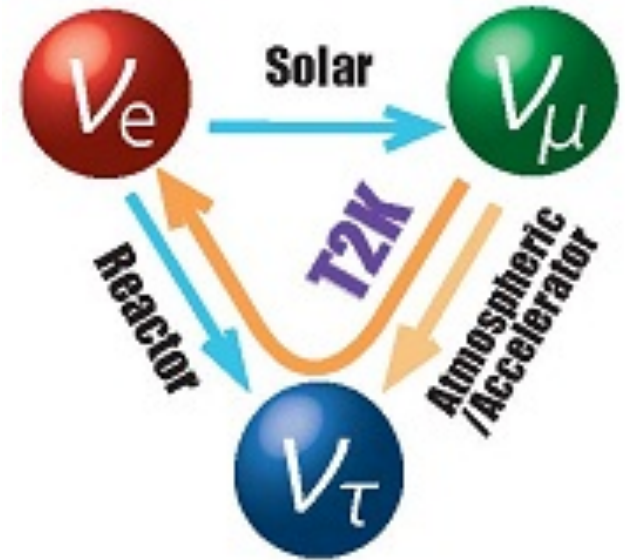
PDG 2020

$$\theta_{23} \simeq 49.2 \pm 1^\circ, \quad \theta_{13} = 8.6 \pm 0.1^\circ, \quad \theta_{12} = 33.4 \pm 0.8^\circ \quad (1\sigma)$$



# Neutrino oscillation

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From atmospheric reactor solar Majorana phase

PDG 2020

$$\theta_{23} \simeq 49.2 \pm 1^\circ, \quad \theta_{13} = 8.6 \pm 0.1^\circ, \quad \theta_{12} = 33.4 \pm 0.8^\circ \quad (1\sigma)$$

$\sin \theta_{23} \simeq 1/\sqrt{2}$ , bi-maximal?

$\sin \theta_{12} \simeq 1/\sqrt{3}$ , tri-maximal?

# $\mu$ - $\tau$ symmetry

Fukuyama and Nishiura 97,  
C. S. Lam 01, E. Ma and M. Raidal 01

bi-maximal 混合を保証する  $Z_2$  symmetry

$$T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad T m_f T = m_f, \quad m_f = \begin{pmatrix} a & b & b \\ c & d & e \\ c & e & d \end{pmatrix}$$

2-3世代の  $45^\circ$  回転 (最大混合) で変換すると、

$$a \sim e \in \mathbb{C}$$

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} a & b & b \\ c & d & e \\ c & e & d \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} = \begin{pmatrix} a & \sqrt{2}b & 0 \\ \sqrt{2}c & d+e & 0 \\ 0 & 0 & d-e \end{pmatrix}$$

残りは1-2回転だけで対角化

もう一方の質量行列が対角的な基底で  $\Rightarrow \theta_{23} = \pm 45^\circ$

# Note: TBMと $K_4 = Z_2 \times Z_2$ 対称性

C.S. Lam, PRD 74, 113004, (2006), PLB 656, 193-198, (2007), PRL 101, 121602 (2008).

2つの離散対称性が厳密なTri-Bi-Maximal 混合を予言

•  $\mu$ - $\tau$  symmetry

2-3 flavorの交換

$$T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad T m_f T = m_f,$$

• magic symmetry

どんな行や列の和も

等しい (魔法陣)

$$S = \begin{pmatrix} -\frac{1}{3} & \frac{2}{3} & \frac{2}{3} \\ \frac{2}{3} & -\frac{1}{3} & \frac{2}{3} \\ \frac{2}{3} & \frac{2}{3} & -\frac{1}{3} \end{pmatrix} \quad S m_f S = m_f,$$

C. S. Lam, Phys. Lett. B 640 (2006) 260, P. F. Harrison and W. G. Scott, Phys. Lett. B 594, 324 (2004)

R. Friedberg and T. D. Lee, HEPNP 30, 591 (2006).

この $K_4 = Z_2 \times Z_2$  対称性 ( $\triangleleft A_4$ ) を課すとTBMで対角化

$$\begin{pmatrix} \sqrt{\frac{2}{3}} & -\frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{6}} \\ \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} c+d-b & b & b \\ b & c & d \\ b & d & c \end{pmatrix} \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} = \begin{pmatrix} c+d-2b & 0 & 0 \\ 0 & c+d+b & 0 \\ 0 & 0 & c-d \end{pmatrix}$$

3 parameters

それぞれの $Z_2$  対称性が固有ベクトルを固定

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線形代数の試験問題を  
無限に生成可能

C. S. Lam, Phys. Lett. B 640 (2006)

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# $\mu$ - $\tau$ reflection symmetry

Harrison & Scott, '02  
Grimus & Lavoura '02

=  $\mu$ - $\tau$  symmetry + 複素共役

$$T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad T m_f^* T = m_f, \quad m_f = \begin{pmatrix} a = a^* & b & b^* \\ c & d & e \\ c^* & e^* & d^* \end{pmatrix}$$

$a \sim e \in \mathbb{C}$

23 最大混合で回すと、

$$U_{BM} m_f U_{BM}^\dagger = \begin{pmatrix} a & \sqrt{2} \operatorname{Re} b & -\sqrt{2} i \operatorname{Im} b \\ \sqrt{2} \operatorname{Re} c & \operatorname{Re}(d+e) & i \operatorname{Im}(e-d) \\ -\sqrt{2} i \operatorname{Im} c & i \operatorname{Im}(-e-d) & \operatorname{Re}(d-e) \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -i \end{pmatrix} m_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & i \end{pmatrix},$$

$m_0$  は一般的な実行列 (9 parameters).

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$m_0$  は一般的な実行列 (9 parameters).

もう一方の質量行列が対角的な基底で  $\Rightarrow \theta_{23} = \pm 45^\circ$   
 $\delta_{CP} = \pm 90^\circ,$

# $\mu$ - $\tau$ reflection symmetry

Harrison & Scott, '02  
Grimus & Lavoura '02

$$U_{BM} m_f U_{BM}^\dagger = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -i \end{pmatrix} m_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & i \end{pmatrix},$$

実行列  $m_0$  は直交行列で対角化

$$R_L m_0 R_R^T = \begin{pmatrix} \pm m_1 & 0 & 0 \\ 0 & \pm m_2 & 0 \\ 0 & 0 & \pm m_3 \end{pmatrix}. \quad \Rightarrow \alpha_{2,3}/2 = 0 \text{ or } \pm\pi/2$$

負の固有値  $\Rightarrow$   
最大マヨラナ位相

$\mu$ - $\tau$  reflection は一般化されたCP対称性のひとつ

a generalized CP (GCP) symmetry

Ecker, Grimus, Konetschny '80,  
Holthausen et al, '12, Feruglio, et al, '12

$$\psi'_i = U_{ij} C \psi_j^*(t, -\mathbf{x}), \quad m'_\psi = U_L m_\psi^* U_R^\dagger,$$



# $\mu$ - $\tau$ (reflection) symmetry

Fukuyama and Nishiura 97,  
C. S. Lam 01, E. Ma and M. Raidal 01

- $\mu$ - $\tau$  symmetry  $Tm_fT = m_f,$

$$T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$m_f = \begin{pmatrix} a & b & b \\ c & d & e \\ c & e & d \end{pmatrix}$$

$$a \sim e \in \mathbb{C}$$

(in the diagonal basis)

$$\Rightarrow \theta_{23} = \pm 45^\circ$$

- $\mu$ - $\tau$  reflection symmetry  $Tm_f^*T = m_f,$

Harrison & Scott, '02

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$$\Rightarrow \theta_{23} = \pm 45^\circ$$

$$\delta_{CP} = \pm 90^\circ,$$

$$\alpha_{2,3}/2 = 0 \text{ or } \pm \pi/2$$

- Reflection symmetry  $Rm_f^*R = m_f,$

M. Yang, '20, submitting to Chin.Phys.C

$$R = \text{diag}(-1, 1, 1),$$

$$m_f = \begin{pmatrix} A & iB & iC \\ iD & F & G \\ iE & H & K \end{pmatrix}$$

$$A \sim K \in \mathbb{R}$$

$$\delta \sim \pm 90^\circ,$$

$$\alpha_{2,3}/2 = 0 \text{ or } \pm \pi/2$$

# 動機：Universal texture

Koide, Nishiura, Matsuda,  
Kikuchi, Fukuyama, '02

すべてのfermionに23対称性と0 textureを課す

$$T\hat{m}_f T = \hat{m}_f, \quad (m_f)_{11} = 0,$$

$$\hat{m}_f = \begin{pmatrix} 0 & A_f & A_f \\ A_f & B_f & C_f \\ A_f & C_f & B_f \end{pmatrix} \quad A_f \sim C_f \in \mathbb{R}$$

質量と混合がだいたい上手くいくが、  
 $\Rightarrow \sin \theta_{13} \simeq 0.05$   
2020現在:  $\sin \theta_{13} \simeq 0.15$

課題： $\theta_{13}$ を大きくしたい

CP位相も含めて統一的textureにしたい

$\Rightarrow \mu$ - $\tau$  reflection symmetryを用いる

# 統一的な $\mu$ - $\tau$ reflection

M. Yang, '20, PLB 806

- 同じ $\mu$ - $\tau$  reflectionを課すと、相殺して $\delta_{\text{CKM}} = 0$ になる
- $\Rightarrow$  別々の $\mu$ - $\tau$  reflection symmetriesが必要

$$T_u(m_u^{BM})^* T_u = m_u^{BM}, \quad T_d(m_d^{BM})^* T_d = m_d^{BM},$$

$$T_u = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad T_d = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & -1 & 0 \end{pmatrix}$$

また、仮定として  $m_f^\dagger = m_f$ ,  $(m_f)_{11} = 0$ ,  $(m_f)_{12} = (m_f)_{13}$

$\rightarrow$  質量固有値とCKM混合はOK

# Four-zero textureとの関係

階層的基底でfour-zero textureになる

$$m_f^{BM} \equiv U_{BM}^\dagger m_f U_{BM},$$

$$m_u = \begin{pmatrix} 0 & iC_u & 0 \\ -iC_u & \tilde{B}_u & B_u \\ 0 & B_u & A_u \end{pmatrix} \quad m_d = \begin{pmatrix} 0 & C_d & 0 \\ C_d & \tilde{B}_d & B_d \\ 0 & B_d & A_d \end{pmatrix} \quad \text{Fritzsch \& Xing, 95}$$

$A \sim K \in \mathbb{R}$

(元々のUniversal textureでは $B_f = 0$ )

近似的に  $12 \times 23$  混合のみで対角化  $\Rightarrow$  CKM行列の良い記述

$$V_{\text{CKM}} = O_u^\dagger \text{diag}(-i, 1, 1) O_d$$
$$= \begin{pmatrix} c_u & s_u & 0 \\ -s_u & c_u & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{-i\delta_{\text{FZ}}} & 0 & 0 \\ 0 & c_q & s_q \\ 0 & -s_q & c_q \end{pmatrix} \begin{pmatrix} c_d & -s_d & 0 \\ s_d & c_d & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$\mu - \tau$  reflection symmetries  $\Rightarrow \delta_{\text{FZ}} = 90^\circ$       観測  $\delta_{\text{FZ}} = 87.9^\circ$

# Four-zero textureとの関係

## 階層的基底でfour-zero textureになる

$$m_f^{BM} \equiv U_{BM}^\dagger m_f U_{BM},$$

Hermite性を課したとき、

$(mf)_{11}=0$ で1つ,  $(mf)_{13}=(mf)_{31}=0$ で1つ

$$m_u = \begin{pmatrix} 0 & iC_u & 0 \\ -iC_u & \tilde{B}_u & B_u \\ 0 & B_u & A_u \end{pmatrix} \quad m_d = \begin{pmatrix} 0 & C_d & 0 \\ C_d & \tilde{B}_d & B_d \\ 0 & B_d & A_d \end{pmatrix} \quad \text{Fritzsch \& Xing, 95}$$

$$A \sim K \in \mathbb{R}$$

(元々のUniversal textureでは $B_f = 0$ )

近似的に  $12 \times 23$  混合のみで対角化  $\Rightarrow$  CKM行列の良い記述

$$\begin{aligned} V_{\text{CKM}} &= O_u^\dagger \text{diag}(-i, 1, 1) O_d \\ &= \begin{pmatrix} c_u & s_u & 0 \\ -s_u & c_u & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{-i\delta_{\text{FZ}}} & 0 & 0 \\ 0 & c_q & s_q \\ 0 & -s_q & c_q \end{pmatrix} \begin{pmatrix} c_d & -s_d & 0 \\ s_d & c_d & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{aligned}$$

$$\mu - \tau \text{ reflection symmetries} \Rightarrow \delta_{\text{FZ}} = 90^\circ \quad \text{観測} \quad \delta_{\text{FZ}} = 87.9^\circ$$

# Four-zero textureとの関係

## 階層的基底でfour-zero textureになる

$$m_f^{BM} \equiv U_{BM}^\dagger m_f U_{BM},$$

Hermite性を課したとき、

$(mf)_{11}=0$ で1つ,  $(mf)_{13}=(mf)_{31}=0$ で1つ

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**Maximal CPV !**

$\mu - \tau$  reflection symmetries  $\Rightarrow \delta_{\text{FZ}} = 90^\circ$       観測  $\delta_{\text{FZ}} = 87.9^\circ$

# Reflection symmetries

M. Yang, '20, Chin. Phys. Cに投稿中

この階層的基底において、 $\mu$ - $\tau$  reflectionは  
対角的なgeneralized CP 対称性になる  $\psi'_i = U_{ij} C \psi_j^*(t, -\mathbf{x})$ ,

Ecker, Grimus, Konetschny '80, Holthausen et al, '12, Feruglio, et al, '12

$$-U_{BM}^* T_u U_{BM}^\dagger = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \equiv R, \quad U_{BM}^* T_d U_{BM}^\dagger = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = 1_3.$$

$$R m_u^* R = m_u,$$

$$m_d^* = m_d.$$

$$m_u = \begin{pmatrix} a_u & \boxed{ib_u} & \boxed{ic_u} \\ \boxed{-ib_u} & d_u & e_u \\ \boxed{-ic_u} & e_u & f_u \end{pmatrix}, \quad m_d = \begin{pmatrix} a_d & b_d & c_d \\ b_d & d_d & e_d \\ c_d & e_d & f_d \end{pmatrix},$$

pure imaginal

片方の第1世代のみに  
CP位相を集中する

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pure imaginal

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$\mu$ - $\tau$  reflection symmetriesから $\mu$ - $\tau$  symmetry を引いたもの

→ reflection symmetries



# Universal four-zero texture

すべてのfermionに同じ対称性とtextureを課す

$$\begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \equiv R, \quad Rm_{u,\nu}^*R = m_{u,\nu}, \quad m_{d,e}^* = m_{d,e}.$$
$$m_u = \begin{pmatrix} 0 & iC_u & 0 \\ -iC_u & \tilde{B}_u & B_u \\ 0 & B_u & A_u \end{pmatrix}, \quad m_d = \begin{pmatrix} 0 & C_d & 0 \\ C_d & \tilde{B}_d & B_d \\ 0 & B_d & A_d \end{pmatrix}$$
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**Concise and  
Elegant !!**

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レプトンは  $4 \times 2$  parameter  $\leftarrow$  3 lepton mass, 2 neutrino mass, 3 mixing

$\Rightarrow m1, \delta CP, \alpha_{2,3}$ を決定

$$U_{\text{MNS}} = V_e^T \begin{pmatrix} -i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} V_\nu P_M,$$

# four-zero textureの正当性 (またあとで)

- YukawaのHermite性

←  $SU(2)_L \times SU(2)_R$  modelでのパリティ対称性

$$\mathcal{L} \ni \bar{\psi}_{Li} Y_{fij} \Phi \psi_{Rj} + \bar{\psi}_{Ri} Y_{fij}^\dagger \Phi^* \psi_{Lj},$$

$$P : \psi_L \leftrightarrow \psi_R, \quad \Phi^* \leftrightarrow \Phi \quad \Rightarrow \quad Y_f = Y_f^\dagger$$

- $(m_f)_{11} = 0$

しばしば離散対称性で実現

第1世代の質量を禁止するカイラル対称性を示唆

- $(m_f)_{12} \neq 0, (m_f)_{13} = 0,$

やや非自明だが、LR modelで正当化可能 (あとで)

# Calculation of $U_{\text{MNS}}$

$\nu$ は階層が弱い→対角化は13混合を無視できない

$$V_{\text{CKM}} \simeq \begin{pmatrix} 1 & \sqrt{\frac{m_u}{m_c}} & 0 \\ -\sqrt{\frac{m_u}{m_c}} & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} -i & 0 & 0 \\ 0 & c_q & s_q \\ 0 & -s_q & c_q \end{pmatrix} \begin{pmatrix} 1 & -\sqrt{\frac{m_d}{m_s}} & 0 \\ \sqrt{\frac{m_d}{m_s}} & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$U_{\text{MNS}} = UP, \quad P \equiv \text{diag}(1, e^{i\alpha_2/2}, e^{i\alpha_3/2}).$$

$$U \simeq \begin{pmatrix} 1 & \sqrt{\frac{m_e}{m_\mu}} & 0 \\ -\sqrt{\frac{m_e}{m_\mu}} & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} -i & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} \\ 0 & 1 & 0 \\ -s_{13} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$U_{e3}$ はO.K. (an input parameter)

$$|U_{e3}| = |s_{13}^{PDG}| \simeq |is_{13} - \sqrt{\frac{m_e}{m_\mu}} c_{13} s_{23}|.$$

# Calculation of $U_{\text{MNS}}$

Jarlskog invariant  $\rightarrow$  Dirac phase

$$J = -\text{Im} [U_{\mu 3} U_{\tau 2} U_{\mu 2}^* U_{\tau 3}^*] \quad (36)$$

$$\begin{aligned} &\simeq \sqrt{m_e/m_\mu} c_{13} c_{23} [-c_{12} s_{12} s_{23}^2 + s_{13} c_{23} s_{23} (c_{12}^2 - s_{12}^2) \\ &+ s_{13}^2 c_{12} s_{12} c_{23}^2] = -0.0130, \end{aligned} \quad (37)$$

$$\sin \delta_{CP} \simeq \sqrt{\frac{m_e}{m_\mu} \frac{c_{13} s_{23}}{s_{13}}}$$

$$\delta_{CP} \simeq 203^\circ$$

This value is rather close to the best fit for the normal hierarchy and in the  $1\sigma$  region  $\delta_{CP}/^\circ = 217_{-28}^{+40}$  [40]. This

[40] Esteban, Gonzalez-Garcia, Hernandez-Cabezudo, Maltoni, and Schwetz, JHEP 01, 106 (2019), 1811.05487.

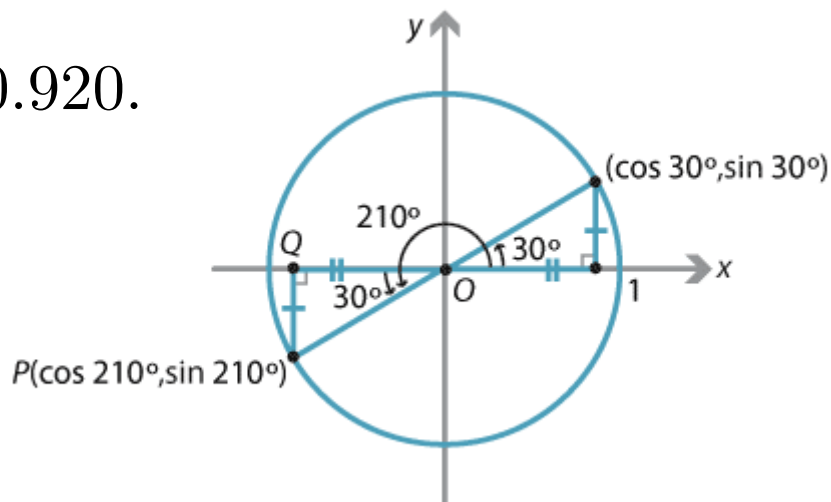
# 保険用ページ

$\delta_{CP}$  は第3象限に存在

$$\cos \delta_{CP} = \frac{|U_{22}^{PDG}|^2 - (s_{12}^{PDG} s_{13}^{PDG} s_{23}^{PDG})^2 - (c_{12}^{PDG} c_{23}^{PDG})^2}{-2s_{12}^{PDG} s_{13}^{PDG} s_{23}^{PDG} c_{12}^{PDG} c_{23}^{PDG}} \quad (40)$$

$$= \frac{|U_{22}|^2(1 - |U_{13}|^2)^2 - |U_{13}|^2|U_{12}|^2|U_{23}|^2 - |U_{11}|^2|U_{33}|^2}{-2|U_{13}||U_{12}||U_{23}||U_{11}||U_{33}|} \quad (41)$$

$$= -0.920. \quad (42)$$



# 物理的予言

ニュートリノ質量行列を再構成できる

$$m_\nu = V_e U_{\text{MNS}} \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix} U_{\text{MNS}}^T V_e^T. \quad (m_e \text{ が 4-zero texture の基底で})$$

0 texture から予言が発生

$$(m_\nu)_{11} = (m_\nu)_{13} = 0 \Rightarrow$$

$$m_1 = \frac{-e^{i\alpha_2} m_2 s_{12}^2 - e^{i\alpha_3} m_3 t_{13}^2}{c_{12}^2}, \quad \simeq 2.5 \text{ [meV]}, \quad (\alpha_2, \alpha_3) = (\pi, 0). \\ \simeq -6.2 \text{ [meV]}, \quad (\alpha_2, \alpha_3) = (0, 0).$$

Double beta decay の effective mass

$$|m_{ee}| = \left| \sum_{i=1}^3 m_i U_{ei}^2 \right| \simeq 1.24 \text{ [meV]} \quad (\alpha_2, \alpha_3) = (\pi, 0) \quad (\text{逆階層では解なし、}) \\ \simeq 0.17 \text{ [meV]} \quad (\alpha_2, \alpha_3) = (0, 0). \quad \text{texture と矛盾)}$$



# 再構成された質量行列の数値

(摂動が悪いかも、なので2~5%くらい誤差あるかも)

Z.-z. Xing & Z.-h. Zhao, NPB 897, '15

$$Y_u \simeq \frac{0.9m_t\sqrt{2}}{v} \begin{pmatrix} 0 & 0.0002i & 0 \\ -0.0002i & \boxed{0.10} & 0.31 \\ 0 & 0.31 & 1 \end{pmatrix} \quad Y_d \simeq \frac{0.9m_b\sqrt{2}}{v} \begin{pmatrix} 0 & 0.005 & 0 \\ 0.005 & \boxed{0.13} & 0.31 \\ 0 & 0.31 & 1 \end{pmatrix}$$

$O(mt)$   $O(mb)$

$m_1 \simeq 2.5[\text{meV}]$ ,  $(\alpha_2, \alpha_3) = (\pi, 0)$ . の解

$$m_{\nu 0} \simeq \begin{pmatrix} 0 & -8.86i & 0 \\ -8.86i & 29.3 & 26.4 \\ 0 & 26.4 & 14.6 \end{pmatrix} [\text{meV}] \quad m_e \simeq \begin{pmatrix} 0 & -7.058 & 0 \\ -7.058 & 107.873 & \boxed{96.12} \\ 0 & 96.12 & \boxed{1740} \end{pmatrix} [\text{MeV}]$$

$O(m_\tau)$

(もう一方の解は $(m_e)_{22} \sim m_\tau$ でやや不自然)

## 素朴には d-e 統一に失敗...??

(エルミート性か $(m_d)_{13} = 0$ を放棄...?)

# シーソー機構

Minkowski '77, Yanagida '79,  
Mohapatra & Senjanovic '80

Four-zero textureとreflection symmetriesはSeesaw不変

⇒  $Y_\nu$ にこれを課せば $M_R$ も従う Nishiura, Matsuda, and Fukuyama '99,

$$Y_\nu = Y_u \simeq \frac{0.9m_t\sqrt{2}}{v} \begin{pmatrix} 0 & 0.0002i & 0 \\ -0.0002i & 0.10 & 0.31 \\ 0 & 0.31 & 1 \end{pmatrix}$$

とすると

$$M_R = \frac{v^2}{2} Y_\nu m_{\nu 0}^{-1} Y_\nu^T \\ = \begin{pmatrix} 0 & -1.08i \times 10^8 & 0 \\ -1.08i \times 10^8 & 1.26 \times 10^{14} & 4.07 \times 10^{14} \\ 0 & 4.07 \times 10^{14} & 1.32 \times 10^{15} \end{pmatrix} [\text{GeV}].$$

$$\Rightarrow RM_R^*R = M_R. \quad (M_{R1}, M_{R2}, M_{R3}) \\ = (2.86 \times 10^6, 3.73 \times 10^9, 1.44 \times 10^{15}) [\text{GeV}],$$

# くりこみに対する安定性

Four-zero texture と reflection sym. は量子補正をほとんど受けない

Fritzsch & Xing, PLB 413, '97  
Xing & Zhao, NPB 897, '15

$$16\pi^2 \frac{dY_u}{dt} = [\alpha_u + C_u^u(Y_u Y_u^\dagger) + C_u^d(Y_d Y_d^\dagger)] Y_u,$$

$$16\pi^2 \frac{dY_d}{dt} = [\alpha_d + C_d^u(Y_u Y_u^\dagger) + C_d^d(Y_d Y_d^\dagger)] Y_d,$$

$$Y_u Y_u^\dagger Y_d = \begin{pmatrix} 1.17 \times 10^{-9}i & 2.34 \times 10^{-12} + 2.56 \times 10^{-7}i & 7.99 \times 10^{-7}i \\ 6.22 \times 10^{-6} & 0.00140 - 1.17 \times 10^{-9}i & 0.00438 \\ 2.00 \times 10^{-5} & 0.00450 - 3.63 \times 10^{-9}i & 0.0141 \end{pmatrix} \\ \simeq \begin{pmatrix} iC_u \tilde{B}_u C_d & iC_u(B_u B_d + \tilde{B}_u \tilde{B}_d) & iC_u(B_u A_d + \tilde{B}_u B_d) \\ (B_u B_u + \tilde{B}_u \tilde{B}_u)C_d & O(B_u A_u B_d) - i\tilde{B}_u C_u C_d & O(B_u A_u A_d) \\ (A_u B_u + B_u \tilde{B}_u)C_d & O(A_u A_u B_d) - iB_u C_u C_d & O(A_u A_u A_d) \end{pmatrix}.$$

第1世代  $C_{u,d} \sim \sqrt{m_{u,d} m_{c,s}}$  が小さい  $\Rightarrow$  the texture と the sym. は独立に安定

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**$\Rightarrow$  高エネルギー(GUT scale)のflavorやCPを受け継いでいる**

# UV theory?

G-J Ding, S. F. King, A. J. Stuart  
JHEP 1013) 006, ArXiv 1307.4212.

これらは別々に破れた residual symmetries

UV theory with the only GCP sym.

Spontaneous CPV

$$\langle \theta_u \rangle = iV_u$$

$$Rm_{u,\nu}^* R = m_{u,\nu},$$

$$m_u = \left( \begin{array}{c|cc} 0 & iC_u & 0 \\ \hline -iC_u & \tilde{B}_u & B_u \\ 0 & B_u & A_u \end{array} \right)$$

$$\langle \theta_d \rangle = V_d$$

$$m_{d,e}^* = m_{d,e}.$$

$$m_d = \left( \begin{array}{c|cc} 0 & C_d & 0 \\ \hline C_d & \tilde{B}_d & B_d \\ 0 & B_d & A_d \end{array} \right)$$

flavored CPV (もしくはSCPV)が、第1世代のみ特別扱い

第1世代の質量を禁止するカイラル対称性の破れに付随??

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$\Rightarrow U(1)_{PQ}$  symmetry?

# Realization of reflection symmetries

Model : 2HDM + 2 flavons  $\theta_{u,d}$  +  $Z_2^{\text{NFC}}$   $\times$   $U(1)_{\text{PQ}}$   $\times$  GCP

- $Z_2^{\text{NFC}}$  : FCNCを禁止 + ポテンシャルを制限
- $U(1)_{\text{PQ}}$  : 第1世代の質量を禁止 +  $\bar{\theta} = 0$
- GCP : 湯川の複素位相を制御

類似の模型がK. Kang and M. Shin,  
Phys. Rev. D 33, 2688 (1986).

|            | $SU(2)_L$ | $U(1)_Y$ | $Z_2^{\text{NFC}}$ | $U(1)_{\text{PQ}}$ | $CP$ |
|------------|-----------|----------|--------------------|--------------------|------|
| $q_{Li}$   | <b>2</b>  | 1/6      | 1                  | -1, 0, 0           | 1    |
| $u_{Ri}$   | <b>1</b>  | 2/3      | 1                  | 1, 0, 0            | 1    |
| $d_{Ri}$   | <b>1</b>  | -1/3     | -1                 | 1, 0, 0            | 1    |
| $l_{Li}$   | <b>2</b>  | -1/2     | 1                  | -1, 0, 0           | 1    |
| $\nu_{Ri}$ | <b>1</b>  | 0        | 1                  | 1, 0, 0            | 1    |
| $e_{Ri}$   | <b>1</b>  | -1       | -1                 | 1, 0, 0            | 1    |
| $H_u$      | <b>2</b>  | -1/2     | 1                  | 0                  | 1    |
| $H_d$      | <b>2</b>  | 1/2      | -1                 | 0                  | 1    |
| $\theta_u$ | <b>1</b>  | 1        | 1                  | -1                 | +i   |
| $\theta_d$ | <b>1</b>  | 1        | -1                 | -1                 | -i   |

$\theta_{u,d}$ と第1世代のみ  
 $U(1)_{\text{PQ}}$  chargeを持つ

$\theta_{u,d}$ のみ  
GCP chargeを持つ

$\theta_{u,d}$ のvevで  
 $U(1)_{\text{PQ}}$  とGCPが破れる



# 1. flavored U(1)<sub>PQ</sub>

F. Wilczek, Phys.Rev.Lett. 49, 1549 (1982),

A. Davidson and K. C. Wali, Phys. Rev. Lett. 48, 11 (1982),

Y. Ahn, Phys.Rev.D 91, 056005 (2015).

$\theta_{u,d}$ と第1世代のみU(1)<sub>PQ</sub> chargeを持つ

湯川行列の電荷

$$\begin{aligned} q_{1L} &\rightarrow e^{-i\alpha} q_{1L}, & u_{1R} &\rightarrow e^{i\alpha} u_{1R}, & d_{1R} &\rightarrow e^{i\alpha} d_{1R}, \\ l_{1L} &\rightarrow e^{-i\alpha} l_{1L}, & \nu_{1R} &\rightarrow e^{i\alpha} \nu_{1R}, & e_{1R} &\rightarrow e^{i\alpha} e_{1R}. \end{aligned} \quad \left( \begin{array}{c|cc} e^{2i\alpha} & e^{i\alpha} & e^{i\alpha} \\ \hline e^{i\alpha} & 1 & 1 \\ e^{i\alpha} & 1 & 1 \end{array} \right)$$

最も一般的な湯川相互作用

$$\begin{aligned} -\mathcal{L} \ni & \bar{q}_L (\tilde{Y}_u^0 + \frac{\theta_u}{\Lambda} \tilde{Y}_u^1 + \frac{\theta_u^2}{\Lambda^2} \tilde{Y}_u^2 + \frac{\theta_d^2}{\Lambda^2} \tilde{Y}_u'^2) u_R H_u \\ & + \bar{q}_L (\tilde{Y}_d^0 + \frac{\theta_d}{\Lambda} \tilde{Y}_d^1 + \frac{\theta_u \theta_d}{\Lambda^2} \tilde{Y}_d^2) d_R H_d + h.c., \end{aligned}$$

$$\tilde{Y}_{u,d}^0 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \tilde{d}_{u,d} & \tilde{c}_{u,d} \\ 0 & \tilde{b}_{u,d} & \tilde{a}_{u,d} \end{pmatrix}, \quad \tilde{Y}_{u,d}^1 = \begin{pmatrix} 0 & \tilde{e}_{u,d} & \tilde{f}_{u,d} \\ \tilde{g}_{u,d} & 0 & 0 \\ \tilde{h}_{u,d} & 0 & 0 \end{pmatrix}$$

flavored axion (flaxion or axiflavor) の一種

Y. Ema, K. Hamaguchi, T. Moroi, and K. Nakayama, JHEP 01, 096 (2017),

L. Calibbi, F. Goertz, D. Redigolo, R. Ziegler, and J. Zupan, Phys. Rev. D 95, 095009 (2017).

## 2. Generalized CP

|            | $Z_2^{\text{NFC}}$ | $U(1)_{\text{PQ}}$ | $CP$ |
|------------|--------------------|--------------------|------|
| $H_u$      | 1                  | 0                  | 1    |
| $H_d$      | -1                 | 0                  | 1    |
| $\theta_u$ | 1                  | -1                 | $+i$ |
| $\theta_d$ | -1                 | -1                 | $-i$ |

$\theta_{u,d}$ のみGCP chargeを持つ

$$\theta_u^* = +i\theta_u, \quad \theta_d^* = -i\theta_d, \quad \phi^* = \phi \quad \text{for other fields}$$

ポテンシャルは $Z_2$ と $U(1)_{\text{PQ}}$ で実(GCP inv.)なものしか許されない

$$V = V^1(H_u, H_d) + V^2(H_{u,d}, \theta_{u,d}) + V^3(\theta_u, \theta_d).$$

(4次項の具体例： $|\theta_u|^2|\theta_d|^2$  or  $\theta_u^*\theta_d\theta_u\theta_d$ )

$\theta_{u,d}$ の真空期待値（実で位相なし）でGCPと $U(1)_{\text{PQ}}$ が破れる

⇒ この基底において、CP位相は湯川の第1世代のみに集中

$$\tilde{Y}_{u,d}^0 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \tilde{d}_{u,d} & \tilde{c}_{u,d} \\ 0 & \tilde{b}_{u,d} & \tilde{a}_{u,d} \end{pmatrix}, \quad \tilde{Y}_{u,d}^1 = \begin{pmatrix} 0 & \tilde{e}_{u,d} & \tilde{f}_{u,d} \\ \tilde{g}_{u,d} & 0 & 0 \\ \tilde{h}_{u,d} & 0 & 0 \end{pmatrix}$$

$$(\tilde{Y}_{u,d}^0)^* = \tilde{Y}_{u,d}^0, \quad \tilde{Y}_u^1 = e^{i\pi/4}|\tilde{Y}_u^1|, \quad \tilde{Y}_d^1 = e^{-i\pi/4}|\tilde{Y}_d^1|.$$

# 結果

$$Y_{u,d} = (\tilde{Y}_{u,d}^0 + \frac{\langle \theta_{u,d} \rangle}{\Lambda} \tilde{Y}_{u,d}^1) = \begin{pmatrix} O(\frac{\langle \theta_{u,d} \rangle^2}{\Lambda^2}) & \tilde{e} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{f} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} \\ \tilde{g} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{d}_{u,d} & \tilde{c}_{u,d} \\ \tilde{h} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{b}_{u,d} & \tilde{a}_{u,d} \end{pmatrix},$$

$$\varphi_u = +\pi/4, \quad \varphi_d = -\pi/4. \quad \begin{array}{l} \text{q,u,dの} \\ \longrightarrow \\ \text{位相変換} \end{array} \quad \varphi_u = \pi/2, \quad \varphi_d = 0.$$

quark質量行列のbest fitから、

$$\frac{\langle \theta_u \rangle}{\Lambda} |\tilde{Y}_u^1| \simeq \frac{\sqrt{2m_u m_c}}{v \sin \beta} \simeq \frac{3 \times 10^{-4}}{\sin \beta}, \quad \frac{\langle \theta_d \rangle}{\Lambda} |\tilde{Y}_d^1| \simeq \frac{\sqrt{2m_d m_s}}{v \cos \beta} \simeq \frac{1 \times 10^{-4}}{\cos \beta},$$

---


$$\frac{\langle \theta_{u,d} \rangle^2}{\Lambda^2} \simeq \frac{10^{-8} (\times \tan^2 \beta)}{|\tilde{Y}_{u,d}^1|^2} \lesssim (y_u, y_d) \simeq \left( \frac{m_u}{v \sin \beta}, \frac{m_d}{v \cos \beta} \right) \simeq (10^{-5}, 10^{-5} \tan \beta).$$

11成分は十分小さい  $\rightarrow$  reflection symmetries と  $(m_f)_{11} = 0$  をみたく

(さらに、(別のCPの基底で)  $\tilde{Y}_u^1 = \tilde{Y}_d^1$  であれば  $(m_f)_{13} = 0$  もみたく)

# 結果

$$Y_{u,d} = (\tilde{Y}_{u,d}^0 + \frac{\langle \theta_{u,d} \rangle}{\Lambda} \tilde{Y}_{u,d}^1) = \begin{pmatrix} O(\frac{\langle \theta_{u,d} \rangle^2}{\Lambda^2}) & \tilde{e} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{f} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} \\ \tilde{g} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{d}_{u,d} & \tilde{c}_{u,d} \\ \tilde{h} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{b}_{u,d} & \tilde{a}_{u,d} \end{pmatrix},$$

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---

The strong CP problem

$$10^{-10} \gtrsim \bar{\theta} = \theta_{\text{QCD}} + \theta_{\text{QFD}} (\equiv \text{Arg Det}[m_u m_d]),$$

**エルミート性やミラー粒子なしに  $\delta_{\text{CKM}}$  と  $\theta_{\text{QFD}}^{\text{tree}} = 0$  が両立**

→ PやCPを用いたstrong CPの解にも適用可能？

# 結果

$$Y_{u,d} = (\tilde{Y}_{u,d}^0 + \frac{\langle \theta_{u,d} \rangle}{\Lambda} \tilde{Y}_{u,d}^1) = \begin{pmatrix} O(\frac{\langle \theta_{u,d} \rangle^2}{\Lambda^2}) & \tilde{e} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{f} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} \\ \tilde{g} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{d}_{u,d} & \tilde{c}_{u,d} \\ \tilde{h} \frac{\langle \theta_{u,d} \rangle}{\Lambda} e^{i\varphi_{u,d}} & \tilde{b}_{u,d} & \tilde{a}_{u,d} \end{pmatrix},$$

$$\varphi_u = +\pi/4, \quad \varphi_d = -\pi/4. \quad \begin{array}{l} \text{q,u,dの} \\ \longrightarrow \\ \text{位相変換} \end{array} \quad \varphi_u = \pi/2, \quad \varphi_d = 0.$$

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今回はflavored  $U(1)_{PQ}$ により  $\bar{\theta}$  はdynamicalに消滅、

$$\Lambda_{\text{GUT}} \simeq 10^{16} \text{ [GeV]}, \Rightarrow \langle \theta_{u,d} \rangle \sim \Lambda_{\text{GUT}} \frac{\sqrt{m_{u,d} m_{c,s}}}{v} \sim 10^{12} \text{ [GeV]}.$$

$\Rightarrow$  現象論的制限と矛盾せず、

Y. Ema, K. Hamaguchi, T. Moroi, and  
K. Nakayama, JHEP 01, 096 (2017),

$$m_a \simeq 10^{-6} \text{ [eV]}, \text{ the dark matter abundance } \Omega_a h^2 \sim 0.2.$$

# まとめ

- Reflection sym. というSMの新しいGCPを提案
  - ⇒ Majorana位相  $\alpha_{2,3}/2 = 0$  or  $\pi/2$
  - ⇒ エルミート性やミラー粒子なしに  $\delta_{\text{CKM}}$  と  $\theta_{\text{QFD}}^{\text{tree}} = 0$  が両立
- + universal four-zero texture
  - ⇒ すべてのfermion 質量と  $V_{\text{CKM}}$  &  $U_{\text{MNS}}$  を再現
  - ⇒  $\delta_{\text{CP}} \doteq 203^\circ$ , 通常質量階層、  $m_1 \doteq 2.5$  or  $6.2$  meVを予言
- この構造はseesaw不変なので、 $Y_\nu$ に課せば $M_R$ も従う
- 第1世代の軽さから、くりこみに耐えうる
- + 対称性の実現のために2HDM+U(1)<sub>PQ</sub> × CPでSSBを考えた
  - ⇒  $f_a \sim M_{\text{GUT}} \sqrt{m_{u,d} m_{c,s}} / v \sim 10^{12} \text{ GeV} \Rightarrow \Omega_a h^2 \sim 0.2$ . axion DM?

$$m_u = \begin{pmatrix} 0 & iC_u & 0 \\ -iC_u & \tilde{B}_u & B_u \\ 0 & B_u & A_u \end{pmatrix} \quad m_d = \begin{pmatrix} 0 & C_d & 0 \\ C_d & \tilde{B}_d & B_d \\ 0 & B_d & A_d \end{pmatrix}$$

$$m_\nu = \begin{pmatrix} 0 & iC_\nu & 0 \\ iC_\nu & \tilde{B}_\nu & B_\nu \\ 0 & B_\nu & A_\nu \end{pmatrix}, \quad m_e = \begin{pmatrix} 0 & C_e & 0 \\ C_e & \tilde{B}_e & B_e \\ 0 & B_e & A_e \end{pmatrix},$$

**That's all. Thank you!**

Back ups



# 位相変換後のreflection symmetries

rephasing of quark fields  $Q = q, u, d$

$$Q' = P_Q^\dagger Q, \quad P_Q = \text{diag}(e^{i\phi_Q}, 1, 1),$$

$$R_{q,u} \equiv P_{q,u} R P_{q,u} = \begin{pmatrix} -e^{2i\phi_{q,u}} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \tilde{R}_{q,d} \equiv P_{q,d} 1_3 P_{q,d} = \begin{pmatrix} +e^{2i\phi_{q,d}} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

$$R_q^\dagger \tilde{m}_u^* R_u = \tilde{m}_u, \quad \tilde{R}_q^\dagger \tilde{m}_d^* \tilde{R}_d = \tilde{m}_d.$$

先ほどの例にするには  $\phi_u = 3\pi/4, \phi_q = -\phi_d = \pi/4$

$$\begin{pmatrix} i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \tilde{m}_u^* \begin{pmatrix} i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \tilde{m}_u, \quad \begin{pmatrix} -i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \tilde{m}_d^* \begin{pmatrix} -i & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \tilde{m}_d,$$

$$\varphi_u = +\pi/4, \quad \varphi_d = -\pi/4.$$

# Global Fit

I. Esteban, M.C.Gonzalez-Garcia, M. Maltoni, T. Schwetz, A. Zhou,  
 JHEP 09 (2020) 178 • e-Print: 2007.14792

$$|U|_{3\sigma}^{\text{with SK-atm}} = \begin{pmatrix} 0.801 \rightarrow 0.845 & 0.513 \rightarrow 0.579 & 0.143 \rightarrow 0.155 \\ 0.234 \rightarrow 0.500 & 0.471 \rightarrow 0.689 & 0.637 \rightarrow 0.776 \\ 0.271 \rightarrow 0.525 & 0.477 \rightarrow 0.694 & 0.613 \rightarrow 0.756 \end{pmatrix}$$

|                          |   | Normal Ordering (best fit)      |                               | Inverted Ordering ( $\Delta\chi^2 = 7.1$ ) |                               |
|--------------------------|---|---------------------------------|-------------------------------|--|-------------------------------|
|                          |   | bfp $\pm 1\sigma$               | $3\sigma$ range               | bfp $\pm 1\sigma$                          | $3\sigma$ range               |
| with SK atmospheric data | $\sin^2 \theta_{12}$                              | $0.304^{+0.012}_{-0.012}$       | $0.269 \rightarrow 0.343$     | $0.304^{+0.013}_{-0.012}$                  | $0.269 \rightarrow 0.343$     |
|                          | $\theta_{12}/^\circ$                              | $33.44^{+0.77}_{-0.74}$         | $31.27 \rightarrow 35.86$     | $33.45^{+0.78}_{-0.75}$                    | $31.27 \rightarrow 35.87$     |
|                          | $\sin^2 \theta_{23}$                              | $0.573^{+0.016}_{-0.020}$       | $0.415 \rightarrow 0.616$     | $0.575^{+0.016}_{-0.019}$                  | $0.419 \rightarrow 0.617$     |
|                          | $\theta_{23}/^\circ$                              | $49.2^{+0.9}_{-1.2}$            | $40.1 \rightarrow 51.7$       | $49.3^{+0.9}_{-1.1}$                       | $40.3 \rightarrow 51.8$       |
|                          | $\sin^2 \theta_{13}$                              | $0.02219^{+0.00062}_{-0.00063}$ | $0.02032 \rightarrow 0.02410$ | $0.02238^{+0.00063}_{-0.00062}$            | $0.02052 \rightarrow 0.02428$ |
|                          | $\theta_{13}/^\circ$                              | $8.57^{+0.12}_{-0.12}$          | $8.20 \rightarrow 8.93$       | $8.60^{+0.12}_{-0.12}$                     | $8.24 \rightarrow 8.96$       |
|                          | $\delta_{\text{CP}}/^\circ$                       | $197^{+27}_{-24}$               | $120 \rightarrow 369$         | $282^{+26}_{-30}$                          | $193 \rightarrow 352$         |
|                          | $\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$    | $7.42^{+0.21}_{-0.20}$          | $6.82 \rightarrow 8.04$       | $7.42^{+0.21}_{-0.20}$                     | $6.82 \rightarrow 8.04$       |
|                          | $\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$ | $+2.517^{+0.026}_{-0.028}$      | $+2.435 \rightarrow +2.598$   | $-2.498^{+0.028}_{-0.028}$                 | $-2.581 \rightarrow -2.414$   |

# Fritzsch-Xing parameterization

H.Fritzsch and Z-z. Xing, *Phys.Lett.B* 413 (1997) 396-404, ArXiv 9707215

Another parameterization of the CKM matrix

$$V = \begin{pmatrix} c_u & s_u & 0 \\ -s_u & c_u & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{-i\varphi} & 0 & 0 \\ 0 & c & s \\ 0 & -s & c \end{pmatrix} \begin{pmatrix} c_d & -s_d & 0 \\ s_d & c_d & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
$$= \begin{pmatrix} s_u s_d c + c_u c_d e^{-i\varphi} & s_u c_d c - c_u s_d e^{-i\varphi} & s_u s \\ c_u s_d c - s_u c_d e^{-i\varphi} & c_u c_d c + s_u s_d e^{-i\varphi} & c_u s \\ -s_d s & -c_d s & c \end{pmatrix},$$

Since the first-generation has tiny mass, only the heavy quark mixing receive quantum corrections

# Fritzsch-Xing parameterization

H.Fritzsch and Z-z. Xing, *Phys.Lett.B* 413 (1997) 396-404, ArXiv 9707215

Another parameterization of the CKM matrix

Almost RGE invariant!

$$\begin{aligned}
 V &= \begin{pmatrix} \boxed{c_u} & \boxed{s_u} & 0 \\ \boxed{-s_u} & \boxed{c_u} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \boxed{e^{-i\varphi}} & 0 & 0 \\ 0 & \boxed{c} & \boxed{s} \\ 0 & \boxed{-s} & \boxed{c} \end{pmatrix} \begin{pmatrix} \boxed{c_d} & \boxed{-s_d} & 0 \\ \boxed{s_d} & \boxed{c_d} & 0 \\ 0 & 0 & 1 \end{pmatrix} \\
 &= \begin{pmatrix} s_u s_d c + c_u c_d e^{-i\varphi} & s_u c_d c - c_u s_d e^{-i\varphi} & s_u s \\ c_u s_d c - s_u c_d e^{-i\varphi} & c_u c_d c + s_u s_d e^{-i\varphi} & c_u s \\ -s_d s & -c_d s & c \end{pmatrix},
 \end{aligned}$$

Since the first-generation has tiny mass, only the heavy quark mixing receive quantum corrections

おまけ

# Note: Predictive GUT

Fritzsch and Xing review, 9912358

universal texture and d-e unification has predictions.

$$M_u = \begin{pmatrix} \mathbf{0} & +ix & \mathbf{0} \\ -ix & y & ry \\ \mathbf{0} & ry & z \end{pmatrix},$$

$$M_d = \begin{pmatrix} \mathbf{0} & x' & \mathbf{0} \\ x' & y' & ry' \\ \mathbf{0} & ry' & z' \end{pmatrix},$$

$$M_e = \begin{pmatrix} \mathbf{0} & x' & \mathbf{0} \\ x' & -3y' & ry' \\ \mathbf{0} & ry' & z' \end{pmatrix},$$

**6 parameters!! @ GUT scale**  
(  $r = \sqrt{2}$  )

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 \end{aligned}$$

Solve  
RGEs  
 $\Rightarrow$   
Georgi-  
Jarlskog  
relation

$$m_d = 3m_e \left( 1 + \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \zeta_{de},$$

$$m_s = \frac{m_\mu}{3} \left( 1 - \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \zeta_{de},$$

$$m_b = m_\tau \frac{\xi_t \xi_b^3}{\xi_\tau^3} \zeta_{de}.$$

$$\xi_i = \exp \left[ -\frac{1}{16\pi^2} \int_0^{\ln(M_X/M_Z)} f_i^2(\chi) d\chi \right]$$

6 parameters!! @ GUT scale  
( $r = \sqrt{2}$ )

renormalized mass @ mZ scale



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 \end{aligned}$$

Solve RGEs  $\Rightarrow$  Georgi-Jarlskog relation

$$m_d = 3m_e \left( 1 + \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \zeta_{de},$$

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$$m_b = m_\tau \frac{\xi_t \xi_b^3}{\xi_\tau^3} \zeta_{de} \cdot \text{RGE factor } \sim 3 \text{ from strong ints.}$$

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$$m_s = \frac{m_\mu}{3} \left( 1 - \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \zeta_{de},$$

$$m_b = m_\tau \frac{\xi_t \xi_b^3}{\xi_\tau^3} \zeta_{de} \cdot \text{RGE factor } \sim 3 \text{ from strong ints.}$$

$$\xi_i = \exp \left[ -\frac{1}{16\pi^2} \int_0^{\ln(M_X/M_Z)} f_i^2(\chi) d\chi \right]$$

**6 parameters!! @ GUT scale**  
( $r = \sqrt{2}$ )

renormalized mass @ mZ scale

Taking  $r^2 = 2$  and  $\tan \beta_{\text{susy}} = 50$  for example, we obtain  $m_d \approx 3.6$  MeV,  $m_s \approx 76$  MeV and  $m_b \approx 3.2$  GeV, essentially in agreement with the results listed in (2.9).

# Note: Predictive GUT

Fritzsch and Xing review, 9912358

universal texture and d-e unification has predictions.

$$\begin{aligned}
 M_u &= \begin{pmatrix} 0 & +ix & 0 \\ -ix & y & ry \\ 0 & ry & z \end{pmatrix}, \\
 M_d &= \begin{pmatrix} 0 & x' & 0 \\ x' & \boxed{y'} & ry' \\ 0 & ry' & z' \end{pmatrix}, \\
 M_e &= \begin{pmatrix} 0 & x' & 0 \\ x' & \boxed{-3y'} & ry' \\ 0 & ry' & z' \end{pmatrix},
 \end{aligned}$$

Solve  
RGEs  
 $\Rightarrow$   
Georgi-  
Jarlskog  
relation

$$m_d = 3m_e \left( 1 + \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \boxed{\zeta_{de}},$$

$$m_s = \frac{m_\mu}{3} \left( 1 - \frac{4r^2}{9} \cdot \frac{m_\mu}{m_\tau} \xi_\tau^3 \right) \boxed{\zeta_{de}},$$

$$m_b = m_\tau \frac{\xi_t \xi_b^3}{\xi_\tau^3} \boxed{\zeta_{de}}. \quad \text{RGE factor } \sim 3 \text{ from strong ints.}$$

$$\xi_i = \exp \left[ -\frac{1}{16\pi^2} \int_0^{\ln(M_X/M_Z)} f_i^2(\chi) d\chi \right]$$

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Unfortunately, this scheme contradicts recent observation of  $V_{ub}$  ...

# 再構成された質量行列の数値

(摂動が悪いかも、なので2~5%くらい誤差あるかも)

Xing & Zhao, NPB 897, '15

$$Y_u \simeq \frac{0.9m_t\sqrt{2}}{v} \begin{pmatrix} 0 & 0.0002i & 0 \\ -0.0002i & \boxed{0.10} & 0.31 \\ 0 & 0.31 & 1 \end{pmatrix} \quad Y_d \simeq \frac{0.9m_b\sqrt{2}}{v} \begin{pmatrix} 0 & 0.005 & 0 \\ 0.005 & \boxed{0.13} & 0.31 \\ 0 & 0.31 & 1 \end{pmatrix}$$

$O(mt)$   $O(mb)$

$m_1 \simeq 2.5[\text{meV}]$ ,  $(\alpha_2, \alpha_3) = (\pi, 0)$ . の解

$$m_{\nu 0} \simeq \begin{pmatrix} 0 & -8.86i & 0 \\ -8.86i & 29.3 & 26.4 \\ 0 & 26.4 & 14.6 \end{pmatrix} [\text{meV}] \quad m_e \simeq \begin{pmatrix} 0 & -7.058 & 0 \\ -7.058 & 107.873 & \boxed{96.12} \\ 0 & 96.12 & \boxed{1740} \end{pmatrix} [\text{MeV}]$$

$O(m_\tau)$

(もう一方の解は $(m_e)_{22} \sim m_\tau$ でやや不自然)

## 素朴には d-e 統一に失敗...??

(エルミート性か $(m_d)_{13} = 0$ を放棄...?)